1	The residual mean circulation in the tropical tropopause layer driven by
2	tropical waves
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13 Abstract. We use latent heating estimates derived from rainfall observations to construct 14 model experiments that isolate equatorial waves forced by tropical convection from mid-15 latitude synoptic-scale waves. These experiments are used to demonstrate that quasi-16 stationary equatorial Rossby waves forced by latent heating drive most of the observed residual mean upwelling across the tropopause transition layer within 15° of the equator. 17 18 The seasonal variation of the equatorial waves and the mean meridional upwelling they 19 cause is examined for two full years from 2006 to 2007. We find that changes in 20 equatorial Rossby wave propagation through seasonally varying mean winds is the 21 primary mechanism for producing an annual variation in the residual mean upwelling. In 22 the tropical tropopause layer, averaged within 15° of the equator and between 90-190 hPa, the annual cycle varies between a maximum upwelling of .4 mm s⁻¹ during Boreal 23 winter and Spring and a minimum of .2 mm s⁻¹ during Boreal summer. This variability 24 25 seems to be due to small changes in the mean wind speed in the tropics. Seasonal 26 variations in latent heating have only a relatively minor effect on seasonal variations in 27 tropical tropopause upwelling. We also find that Kelvin waves drive a small downward 28 component of the total circulation over the equator that may be modulated by the quasi-29 biennial oscillation.

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31 **1. Introduction**

Recent variations in stratospheric water vapor with unknown causes have been implicated in global decadal-scale surface temperature variations (Solomon et al., 2010). Most notably, a sudden decrease in stratospheric water vapor in 2001 coincided with a cooling in tropical tropopause temperatures observed in radiosondes (Randel and Wu, 2006), and the change persisted and only slowly recovered through the first decade of the 21st century. These decadal-scale changes to both temperature and water vapor suggest that there was an increase in tropical upwelling across the tropical tropopause layer (TTL). In

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39 this paper we define the TTL to be the layer between 90-190 hPa (13-17 km). The 40 tropical upwelling circulation is driven by waves and deep convection and stratospheric 41 water is further controlled by tropopause temperatures. Hence it is natural to suspect that 42 there is an increase in wave driving, either through enhanced wave forcing or through 43 changes in the pattern of wave propagation that would lead to enhanced wave-flux 44 divergence. Although extratropical wave pumping clearly influences the zonal mean 45 tropical upwelling in the lower stratosphere (Holton et al., 1995), several recent studies 46 have found that global-scale equatorial waves, forced by convection, play an important 47 role in modulating tropical tropopause temperatures and upwelling in the tropopause 48 layer that control stratospheric water vapor (Boehm and Lee, 2003; Kerr-Munslow and 49 Norton, 2006; Norton, 2006; Ryu and Lee, 2010; Grise and Thompson, 2013).

50 Tropical upwelling across the TTL is part of a wave driven (Brewer-Dobson) 51 circulation, and is the region where dehydration of air entering the stratosphere occurs. 52 In this paper we shall quantify mean upwelling in terms of the residual mean vertical 53 velocity. Plumb and Elusczkiewicz (1999) showed that the steady component of the 54 residual mean vertical velocity at the equatorial TTL can only be driven by EP-flux 55 convergence within 20° of the equator in the lower stratosphere. However, Randel et al. 56 (2002) found that transient mid-latitude EP-flux divergence from extra-tropical waves 57 can drive transient upwelling at the equator. The waves that contribute to the TTL 58 upwelling may include synoptic scale transient eddies, extratropical stationary planetary 59 waves (Scott, 2002), gravity waves (Garcia and Randel, 2008; Li et al., 2008; 60 McLandress and Shepherd, 2009; Calvo et al., 2010; Sigmond and Scinocca; 2010) and 61 tropical waves forced by convective heating. It is important to determine the role each of 62 these waves may play in the tropical upwelling so that the causes of variability and long-63 term trends may be more easily identified. It is also important to distinguish variability in 64 forcing from variability that arises from changes in wave propagation due to the mean 65 state variability (Garcia and Randel, 2008; Calvo et al., 2009). Tropical upwelling rates

have been obtained from several studies. Schoeberl et al. (2008) used MLS and HALOE water vapor measurements to determine upwelling rates of .4-.5 mm/s at 17 km (~100 hPa). Dima and Wallace (2007) used ERA-40 reanalysis to determine an annual mean vertical velocity of .55 mm/s. Abalos et al. (2012) derive upwelling rates by several methods using ERA-interim data, which show an annual cycle with amplitude of approximately a factor of two with a minimum in July of ~.25 mm/s and maximum in January of ~.5mm/s

73 It is uncertain which wave type is primarily responsible for driving tropical upwelling. Boehm and Lee (2003) proposed that tropical waves may be important. Kerr-74 75 Munslow and Norton (2006) used ECMWF reanalysis (ERA-15) at 90hPa and deduced 76 that tropical waves made a significant contribution to tropical upwelling. Randel et al. (2008) found that there is an equal contribution to upwelling from both tropical and extra-77 78 tropical waves in ERA-40, and that both tropical and extratropical waves can contribute 79 to the seasonal cycle. They speculate that in the case of tropical waves, this is due to 80 changes in the seasonal pattern of stationary heating. The recent review by Randel and 81 Jensen (2013) mentions that the mechanisms driving the annual cycle in TTL 82 temperatures and upwelling remain to be clarified. Grise and Thompson (2013) found 83 evidence for tropical wave driving of week-to-week variability in TTL upwelling, but 84 results for longer timescales were inconclusive. Deckert and Dameris (2008) found that 85 higher SSTs in the CCM E39/C amplify deep convection, which is correlated in the 86 model with anomalous production of quasi-stationary equatorial waves. (We define 87 quasi-stationary waves here as stationary waves plus waves with low frequency that 88 travel westward relative to the mean flow.) Anomalous EP-flux divergence associated 89 with the quasi-stationary wave anomaly was shown to drive an anomalous low-latitude 90 Brewer-Dobson cell that was responsible for an increase in ozone-poor air of the model 91 lower stratosphere.

92 We attempt to quantify the upwelling driven by equatorial waves, particularly 93 equatorial Rossby waves. Our model simulations are forced using latent heating 94 estimates derived from TRMM rainfall data. Our experiments are designed to remove as 95 much as possible the effects that extra-tropical synoptic-scale waves have on tropical 96 upwelling so that the effects of tropical waves forced by latent heating can be quantified 97 These experiments show that the observed annual mean velocity and separately. 98 temperature field pattern can be reproduced using wave forcing from latent heating alone. 99 The tropical upwelling driven by the quasi-stationary waves in our experiments closely 100 match the upwelling derived from observations. Our experiments further show that the 101 annual cycle in upwelling is not caused by seasonal changes in equatorial wave sources. 102 Instead, the annual cycle in upwelling results primarily from seasonal changes in wave 103 propagation caused by the annual cycle in tropical zonal-mean winds. We shall also 104 examine the separate contributions to upwelling from synoptic, stationary, eastward, and 105 westward waves, and examine mechanisms, such as the quasi-biennial oscillation (QBO), 106 that may give rise to interannual variability in upwelling.

107 **2. Model description**

108 a) Transformed Eulerian-Mean equations

109 The focus of this paper is on the wave-driven vertical transport velocity across the TTL. 110 This velocity is estimated using the residual mean meridional circulation (RMMC), 111 which is part of the Transformed Eulerian-Mean (TEM) formalism. The RMMC 112 velocities are good approximations to Lagrangian-Mean transport velocities when the 113 waves have small amplitude, are steady, and adiabatic. These conditions are reasonable 114 to assume, but it is beyond the scope of this paper to check their validity. A good 115 reference for the TEM formalism and its use for estimating transport is provided by 116 Andrews et al. (1987), Chapters 3 and 9.

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117 The TEM equations have the same form as the Eulerian-Mean equations except 118 that eddy fluxes have been combined into mean meridional circulation components and 119 forcing terms. In addition to thermal wind balance (not displayed), the zonal momentum, 120 thermal energy, and continuity equations express the balance between the RMMC 121 (\bar{v}^*, \bar{w}^*) , the time tendency of the zonal mean zonal wind and temperature, and various 122 forcing mechanisms due to eddies, dissipation, and heating:

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$$\overline{u}_{t} + \overline{v}^{*} \Big[(a\cos\phi)^{-1} (\overline{u}\cos\phi)_{\phi} - f \Big] + \overline{w}^{*} \overline{u}_{z} = (\rho_{0}a\cos\phi)^{-1} \nabla \cdot \mathbf{F} + X; \\
\overline{T}_{t} + a^{-1} \overline{v}^{*} \overline{T}_{\phi} + \overline{w}^{*} R^{-1} H N^{2} = Q; \\
(a\cos\phi)^{-1} (\overline{v}^{*}\cos\phi)_{\phi} + \rho_{0}^{-1} (\rho_{0}\overline{w}^{*})_{z} = 0.$$
(1)

124 The RMMC may be expressed in terms of the Eulerian mean and wave fields, but we will 125 not need the formulae here. The RMMC is calculated instead as a part of the solution to 126 (1). The terms in the equations are zonal wind u, temperature T, Coriolis parameter f, 127 scale height H (7 km), gas constant R, density ρ , earth radius a, and buoyancy 128 frequency N. Zonal mean is denoted with an over-bar. Derivatives with respect to time 129 t, latitude ϕ , and log-pressure altitude z are denoted with a subscript. These equations 130 are forced by the divergence of the Eliassen-Palm (EP) flux vector:

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$$F^{(\phi)} = \rho_{0} a \cos \phi \left(\overline{u}_{z} \left(\overline{v'T'} \right) R(HN^{2})^{-1} - \left(\overline{v'u'} \right) \right),$$

$$F^{(z)} = \rho_{0} a \cos \phi \left\{ \left[f - (a \cos \phi)^{-1} (\overline{u} \cos \phi)_{\phi} \right] \left(\overline{v'T'} \right) R(HN^{2})^{-1} - \left(\overline{w'u'} \right) \right\}$$
(2)

132 Contributions to the heating term Q include radiative and wave heating/cooling.
133 Parameterized dissipative mechanical forcing is denoted by X.

The TEM equations describe how upwelling in the TTL is produced by EP-flux divergence. We will show that he typical wave forcing pattern consists of two subtropical regions where $\nabla \cdot \mathbf{F} < 0$. This is mainly balanced by the Coriolis term on the left side of the zonal momentum equation. In a nearly steady state solution, the timetendency of zonal wind is essentially zero except very close to the equator. Since the Coriolis parameter changes sign across the equator, the TEM meridional wind must also change sign across the equator. By the continuity equation, the resulting meridional gradient in the meridional wind is balanced by a negative vertical gradient in upwelling, hence the peak upwelling lies below the centers of maximum EP-flux convergence. The thermal energy equation says that in a nearly steady state this upwelling must be balanced by radiative heating.

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146 b) Primitive Equation Model

147 All calculations are performed with a time dependent primitive equation model. The 148 model fields are represented by a spherical harmonic expansion with triangular truncation 149 T40. A higher horizontal resolution has no effect on the wave spectrum due to limits on 150 the frequencies resolved in the forcing (Ortland et al. 2011). The vertical log-pressure 151 altitude grid extends from the surface to 60 km with 750 m spacing. Rayleigh friction 152 above 40 km acts on the wave components to prevent wave reflection from the upper boundary. Newtonian cooling with a rate of .05 d^{-1} above the tropopause damps the 153 154 waves and relaxes the zonal mean temperatures of the model to climatological values 155 described below. Waves in the model are forced by heating derived from gridded global 156 TRMM rainfall rate data sampled every three hours (Huffman et al. 2007). The heating 157 is also computed at 3-hourly intervals and interpolated to intermediate model time steps. 158 The heating is tapered to zero poleward of 30°, so the wave response can be interpreted to 159 consist primarily of tropical waves. The vertical heating profiles include both convective 160 and stratiform rain contributions. The zonal mean heating is removed from the forcing 161 introduced into the model since we are not interested in examining the Hadley circulation 162 and the influence it has on zonal mean winds and temperatures. A full description of how 163 the latent heating is derived from TRMM is given in Ryu et al (2011). The Ryu et al. 164 (2011) study examines transient tropical waves in the stratosphere forced by latent heating derived from TRMM in this way, and provides further details on the model and wave spectra. The primary focus of the present paper is on the seasonally-varying model response in the upper troposphere and lower stratosphere to the waves forced by this 3hourly varying heating, but the waves of primary interest for this study are the quasistationary waves.

170 Monthly and zonally averaged zonal winds and temperatures derived from 171 observations are used to define the zonal mean background for the model. The observed 172 values (OBS) were obtained from NCEP reanalysis (Kalnay et al., 1996). The NCEP 173 monthly zonal mean zonal wind values in the troposphere and lower stratosphere were 174 compared to values obtained from Modern-Era Retrospective Analysis (MERRA) and differences were found to be less than 1 ms⁻¹ in the upper troposphere and lower 175 176 stratosphere. These monthly averages were further modified (MOD) as described below. 177 The MOD values are assigned to the midpoint of each model month, and the value used 178 in the model is obtained by linear interpolation to the current model time step. The model zonal mean winds are relaxed to these interpolated values at a rate of .1 d⁻¹ 179 throughout the domain outside the tropics and at a rate of .5 d^{-1} within 30° of the equator. 180 This relaxation rate was required to keep the model zonal mean within 1 ms⁻¹ of 181 182 observations in the tropics. Since the zonal and monthly mean subtropical tropospheric 183 jets are baroclinically unstable, synoptic scale Rossby waves are spontaneously generated 184 in the model when the observed winds and temperatures are used. In order to examine 185 only the waves generated by tropical heating, the MOD values were designed to suppress 186 the spontaneous generation of synoptic waves. The MOD background and TRMM 187 heating from January 2005 through December 2007 is used to produce a multi-year 188 control run (CTRL). The first year is not examined here, and is used for a ramp-up.

189 To construct the MOD background that suppresses synoptic waves, the model is 190 initialized with the OBS values of wind and temperatures for each month plus a random 191 wave disturbance. It is then allowed to run freely for 90 days with no TRMM heating

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192 and with no relaxation to the zonal mean, while synoptic scale waves develop and modify 193 the mean flow into a stable configuration. This duration is long enough for the transient 194 synoptic waves to complete their modification, die off due to the Newtonian cooling and 195 the sponge, and produce a steady zonal mean 'relaxed' state (RLX) plus weak quasi-196 stationary waves. The MOD background is then obtained by merging the RLX and OBS 197 zonal mean fields. The merge retains OBS within 20 degrees of the equator, with a 198 smooth transition to RLX poleward of 35 degrees. Since most of the baroclinic 199 instability arises at mid-latitudes, this modification retains a state that is nearly stable to 200 baroclinic waves, and also retains the tropical zonal mean background state, which is 201 important for producing realistic propagation of the tropical waves forced in the model. 202 The MOD and OBS backgrounds are identical to each other out to 20° and within a few 203 m/s of each other out to about 50° from the equator. The main differences are poleward 204 of 50°. The jet core of MOD is poleward of OBS and has a larger magnitude. We do not 205 expect the equatorially confined tropical waves to be sensitive to the winds poleward 206 of 50°.

207 Two experiments were run to test that the MOD state suppresses the baroclinic 208 waves. In both, the model was initialized with a random disturbance and no tropical 209 waves are forced. However, unlike the simulation that produced the RLX state, the 210 model was constrained to the zonal mean zonal wind with relaxation rates used in CTRL. 211 In the first 'synoptic run', the zonal mean background is relaxed to OBS, while in the 212 second 'modified synoptic run' the mean background is relaxed to the MOD state. Both 213 simulations generate synoptic scale waves that persist for a 3-year duration. However, in 214 the modified case the synoptic waves are smaller by a factor of 5. The amplitude of 215 synoptic waves in this second run is also negligible compared to the tropical waves 216 generated in CTRL, which is the result we are trying to achieve.

217 We will also look at simulations where the background zonal mean zonal wind is 218 relaxed to zero (UBR0) using the same relaxation rates as in CTRL. The zonal mean

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temperatures in UBR0 are set equal to a climatological vertical profile obtained from observations by averaging between 20°N to 20°S. Using zero background winds rather than climatological zonal mean zonal winds also insures that synoptic scale waves do not form on the subtropical jets. Comparing UBR0 and CTRL simulations will enable us to determine the effect that the zonal mean winds have on the wave propagation.

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225 c) TEM model

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226 Our goal is to calculate the component of the RMMC that is forced by the tropical 227 waves. The zonal mean meridional circulation of the CTRL simulation cannot be used 228 for this purpose because the relaxation of the model zonal mean in the CTRL and UBR0 229 simulations severely suppresses that circulation. What we do instead is to compute the 230 EP flux and divergence from the CTRL and UBR0 simulations, then force a zonal mean 231 version of the primitive equation model with this EP-flux divergence. Forced in this way, 232 but without the same relaxation to the zonal mean as in the wave model, the zonally 233 symmetric version of the primitive equation model solves the TEM equations (1) and the 234 meridional and vertical wind fields can be interpreted as the RMMC. The TEM model is 235 run separately for each month. The RMMC becomes steady in about 30 days. This 236 results in an estimate of the RMMC for each month of the 2-year simulation. Averages 237 of three consecutive months give the seasonal estimates shown below.

The zonal momentum forcing in the TEM model is determined by first calculating the zonal wavenumber-frequency spectra of all fields in the control run by using monthly stretches of the model output. The spectral components of each field are then used to calculate the EP-flux divergence spectrum. Following sections describe results when the model is forced with the total spectrum, or the sum of eastward, westward, or stationary parts of this spectrum. The TEM model also retains the Newtonian cooling used in the full wave model simulations. This is represented by the heating term Q in equations (1). It was found that wave heating has negligible effect on the upwelling in the TTL and hence is not included as a forcing term of the TEM model.

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247 The TEM model is initialized with the monthly and zonal mean zonal winds from 248 the 3D simulations, however we found that the RMMC calculated in the model did not 249 depend on these background winds. Although the RMMC becomes steady in about 30 250 days, the zonal wind in the TEM model is not steady, because important tropospheric 251 forcing mechanisms are absent in the simplified model. To help control these tendencies 252 in a way that does not affect the RMMC upwelling in the TTL, the zonal wind is relaxed to the initial state only below 10 km, at a rate of .1 d⁻¹. This relaxation forcing is 253 254 represented by the X term in equations (1). However, the fact that the zonal wind in the 255 TEM model slowly evolves does not affect the RMMC results. Despite continuous 256 evolving zonal winds, the RMMC remains steady for as long as 90 days or more.

257 To determine that the TEM model produces a reasonable estimate of residual 258 mean upwelling, the balance of terms in the TEM equations was examined. It was found 259 that the primary balance in the TEM zonal momentum equation is between the EP-flux 260 divergence and the Coriolis term. The advection terms are negligible within 30° of the 261 equator. Hence the structure of the zonal mean winds in the TEM model has no effect 262 on the RMMC in the tropics. The relaxation forcing X had only a small contribution 263 below 10 km. The zonal wind tendency is also a negligible part of the balance in most 264 cases except in those months where there is some non-zero EP-flux divergence on the 265 equator. However, this tendency also does not affect the residual mean upwelling – it 266 balances a component of the residual mean meridional wind that is symmetric about the 267 equator. This symmetric component of the meridional wind has very small horizontal 268 gradient, hence does not contribute to upwelling through the continuity equation. This 269 was verified by running the TEM model forced by EP-flux divergence that was modified 270 to be zero within 5° of the equator. These test simulations had the same residual mean 271 upwelling as produced by unmodified EP-flux divergence, and also had negligible zonal

wind tendency in the tropics. The balance in the TEM thermal energy equation for all months is between the adiabatic term and Newtonian relaxation and the temperature tendency in the tropics is negligible. Hence the TTL upwelling that we derive is an accurate estimate of the TEM response to tropical wave EP-flux convergence isolated from extratropical wave influences.

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278 **3. Residual mean circulation**

279 a) Heating and mean response

280 Since we will be looking at seasonal variations in the RMMC, it is worth looking first at 281 the seasonal variations in the heating pattern. The latent heating derived from TRMM is 282 averaged over the 2007 seasons and shown in Figure 1. Large regions of heating occur 283 over Africa, Southeast Asia, and South America. Heating amplifies in the summer 284 hemisphere, with the largest magnitude occurring in the DJF season. A band of heating 285 north of the equator is associated with the intertropical convergence zone. This band 286 migrates in latitude with season. The variation in the heating configuration will be seen 287 to produce a seasonal variation in the wave fluxes that is not readily apparent in the EP-288 flux divergences or the RMMC response.

289 We first demonstrate that our simulation produces a realistic stationary wave 290 structure in the tropics in response to the TRMM heating by comparing our simulation to 291 observations presented by Dima and Wallace (2007). Figure 2 shows the results from 292 the control run averaged over the year 2007. The top two panels show the perturbation 293 temperature field with color-shaded contours and the wind field in the plane of the figure 294 with vectors, while the bottom panel shows actual temperatures. The vertical wind 295 components have been scaled by a factor of 1000 so that they are visible. The TTL in 296 this and other figures is marked with red boundary lines. The top panel of Fig. 2 is obtained by averaging over all latitudes within 10° of the equator. 297 Positive vertical 298 velocities mainly occur over deep convective cloudy regions in the western Pacific, while 299 down-welling occurs over clear sky regions. The coldest temperatures occur within the 300 TTL over the upwelling. Our simple model produces a pattern of wind velocities and 301 temperature anomalies above 450 hPa that is very similar in structure and magnitude to 302 that found by Dima and Wallace (2007) in ERA-40 data (their Fig. 5). The bottom panel 303 of Fig. 2 shows temperatures at 100 hPa averaged over 2007 as function of longitude and 304 latitude. It illustrates the western Pacific cold pool, centered between 130°-190° 305 longitude over the equator, with lobes extending westward in the subtropics. In our 306 model this is produced exclusively as part of the stationary Rossby wave pattern. It 307 agrees very well with the ERA-40 data for roughly the same period shown by Virts et al. 308 (2010) (their Fig. 3). We may conclude that tropical waves driven by the spatial 309 variability in the latent heating are primarily responsible for driving observed 310 longitudinal variations in the circulation in the TTL. The structure of the temperatures is 311 also similar in other years of our run.

312 The middle panel of Fig. 2 shows the wave structure in the longitude-latitude 313 plane at 260 hPa (10.5 km) averaged over 2007. The temperature anomalies display the 314 structure of a superposition of several Rossby waves (e.g. Gill, 1980). The Rossby wave 315 symmetric modes represent the response to the component of the heating symmetric 316 about the equator. They have a primary peak away from the equator and a local 317 minimum at the equator, even if the peak heating occurs on the equator. Note that the 318 temperature anomaly is primarily asymmetric about the equator just west of the 319 Greenwich meridian, while it is symmetric elsewhere. This annual mean pattern is also 320 persistent from year to year in the model simulation. These patterns are very similar to 321 the streamfunction response to annually averaged latent heating shown in Schumacher et 322 al. (2004) (their Fig. 5).

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324 b) Seasonal structure

325 We next examine the EP-flux from the full tropical wave spectrum and its 326 divergence, averaged over the four seasons of 2007. Figure 3 shows the result for the 327 UBR0 simulation in the panels on the left and the results for CTRL simulation on the 328 right. The EP-flux is shown by scaled vectors and the flux divergence is shown with 329 colored contours. In all panels of Fig. 3 we see two regions of EP-flux convergence on 330 either side of the equator centered close to 100 hPa. It is this convergence that drives the 331 RMMC upwelling through the TTL over the equator, as explained in Section 2a. The EP-332 flux convergence is more asymmetric about the equator during solstice, with larger 333 magnitude in the summer hemisphere in which the peak heating occurs, as shown in 334 Figure 1.

335 One can see several differences in EP-flux convergence near the tropopause in the 336 two simulations that is caused by the mean zonal winds. First, in the CTRL simulation 337 the convergence is more confined equatorward of 30° than in the UBR0 simulation. This 338 is probably due to the tropical confinement of the stationary Rossby waves by the strong 339 westerlies in the subtropical jets. Second, the CTRL convergence is stronger in general, 340 but particularly so in DJF. As we shall see below, this is the cause of the seasonal 341 variability of the RMMC. Third, for the JJA and SON seasons, the convergence pattern 342 is lower in CTRL than in UBRO, and there are strong regions of EP-flux divergence over 343 the equator. This is possibly related to the arrival of the westward phase of the quasi-344 biennial oscillation (QBO) in the lower stratosphere at this time. For the DJA and MAM 345 seasons there is a weak remnant of the eastward phase of the QBO just above the 346 tropopause. The westward QBO winds confine the Rossby waves slightly lower in the 347 TTL and allow the propagation of Kelvin waves into the stratosphere. The dissipation of 348 Kelvin waves produces the weak positive EP-flux divergence seen directly over the 349 equator in the TTL in the last two seasons of 2007. These features will be further 350 discussed in the following subsections.

351 The EP-flux divergences shown in Fig. 3 are used to force the zonally symmetric 352 TEM model as described in Section 2c. The mean meridional and vertical winds 353 calculated in this way give the RMMC, which provide an estimate for the zonal mean 354 transport velocity. The RMMC's that result from the EP-flux divergences calculated 355 from the UBR0 and CTRL simulations are shown in the left and right columns, 356 respectively, of Figure 4. These Figures use scaled vectors for the RMMC and color 357 contours showing the residual mean vertical velocity. The EP-flux convergence on either 358 side of the equator in the TTL gives rise to two circulation cells with rising motion over 359 the equator and sinking motion in the subtropics.

360 The circulation pattern for each season shown for the UBR0 simulation in Fig. 4 361 is nearly symmetric about the equator, and except for DJF07, the vertical velocity is maximum over the equator. The deviation from symmetry about the equator in the 362 363 circulation results from the asymmetry of the heating about the equator. The local 364 minimum in the vertical velocity over the equator for DJF07 will be discussed below. 365 Comparing the left and right columns, we note that the inclusion of background winds 366 produces considerable asymmetry in the circulation, and both weakens the circulation 367 pattern somewhat and lowers that altitude of maximum vertical velocity by about 1 km. 368 The sensitivity of the RMMC to the structure of the mean winds suggests a mechanism 369 for controlling the seasonal and interannual variation in the RMMC that will be examined 370 further below. In our model the seasonally averaged value of the residual mean vertical 371 velocity at 111 hPa over the equator is between .3 and .6 mm s⁻¹. This is roughly 372 consistent with observational estimates (Schoeberl et al. 2008; Dima and Wallace 2007; 373 Randel et al., 2008; Abalos et al., 2012). The strength of the circulation varies throughout 374 the year.

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376 *c)* Variability

The seasonal variation in the upwelling through the TTL will depend on the variability of the heating, as shown in Fig. 1, and on the variability of the zonal mean basic state through which the tropical waves propagate. We shall explore which mechanisms are responsible for producing variability in the upwelling by making a comparison between the CTRL and UBR0.

382 The observed tropical zonal mean winds for 2006-2007 are shown in Figure 5. 383 The top panel shows monthly and latitudinally averaged zonal mean zonal wind up to 30 384 km. The bottom panel shows the zonal mean zonal winds averaged over the TTL. Both 385 panels show weak, but significant, variability in the upper troposphere. Weak easterly wind speeds on the order of 5-10 ms⁻¹ develop from June through November, while weak 386 westerly wind speeds less than 5 ms⁻¹ prevail at other times. Manabe et al. (1974) first 387 388 described this variability with a general circulation model. The tropospheric jets 389 strengthen and extend into the subtropics of the winter hemisphere, and weaken and 390 retreat into the summer hemisphere. In the stratosphere, the westerly phase of the quasi-391 biennial oscillation (QBO) descends to the top of the TTL between October 2006 and 392 July 2007. Easterly winds lie above the TTL at all other times.

393 The variation of the vertical component of EP-flux, shown in Figure 6, indicates 394 how the background winds affect wave propagation. The left panels of the figure show 395 results of the CTRL simulation and the right panels show results from the UBR0 396 simulations. The top panels show the time evolution of the vertical profiles of vertical 397 EP-flux at the equator and the bottom panels show latitudinal cross sections in the TTL. 398 In the CTRL simulation we see that the upward flux is essentially zero in the tropics 399 during northern summer. When the background winds are set to zero there is still 400 significant wave propagation in the tropics during this period, so we can conclude that the 401 modulation of wave flux is due to background winds and not due to variability in the 402 latent heating wave source. Notice that the reduced vertical wave flux occurs at the time 403 when weak zonal mean easterly winds prevail in the tropical troposphere. These tropical 404 easterlies inhibit the propagation of Rossby waves, as we might expect. The seasonal 405 variation of the latent heating does produce modulation of wave flux in the subtropics, 406 with maximum flux in the summer hemisphere, as seen by comparing the bottom panels 407 of Fig. 6. Apparently the extension of tropical easterlies into the northern summer 408 subtropics also acts to reduce the subtropical wave flux.

409 The modulation of wave flux due to variations in heating and wave propagation 410 leads to a modulation of the EP-flux divergence, and hence a modulation of the RMMC. 411 The variability of the EP-flux divergence in our model simulations is shown in Figure 7. 412 The pattern of variability in the vertical profiles of the divergence in the tropics (top 413 panels) closely resembles the variability in the flux profiles. Once again, the mean winds 414 produce a regular seasonal variability in the EP-flux divergence which minimizes 415 between June and October. The latitude structure of the divergence (bottom panels) also 416 has a strong seasonal variation when observed background winds are used in the model, 417 but the variability is much weaker when the background winds are zero.

418 The modulation of EP flux divergence leads to the modulation of the RMMC 419 shown in Figure 8. There is considerable variability in the upwelling that can be 420 attributed to heating alone (right panels for the UBR0 simulation). Although there is a 421 seasonal pattern evident in the vertical component of EP flux in the UBR0 simulation 422 (Fig. 6, right panels), that pattern is not evident in the upwelling. There is considerable 423 variability in the vertical velocity in this two year simulation, but it is not clear whether 424 this variability is part of a regular seasonal pattern. However, there is a clear seasonal 425 variation in the upwelling for the CTRL simulation that use observed mean winds (left 426 The weak tropical easterlies, as pointed out above, suppress Rossby wave panel). 427 propagation (and enhance Kelvin wave propagation) from June to October and hence this 428 is also the period when the upwelling is weakest. The maximum upwelling during Boreal winter and Spring, averaging around .8 mm s⁻¹ is a factor of two larger than the minimum 429

430 of around .4 mm s⁻¹ from June to October. The mean winds also significantly alter the 431 latitude structure of the upwelling. The peak upwelling through the TTL occurs in the 432 subtropics of the summer hemisphere. There also appear to be enhancements of the 433 upwelling due to anomalies in the heating structure. For example, upwelling in both 434 simulations is strongest during January 2007.

435

436 d) Spectral decomposition

437 As described in Section 2, the wave fields over each of the seasonal 3-month intervals in 438 the CTRL and UBR0 multi-year model simulations were decomposed into frequency 439 components using the fast Fourier transform, from which the EP-flux divergence for each 440 separate frequency was calculated. In the previous subsection the response to the sum of 441 all frequencies was examined. In this subsection we examine the response to forcing the 442 TEM model with the separate contributions from stationary, eastward, or westward 443 traveling waves. In this spectral decomposition over the 3-month interval, 'stationary' 444 waves have periods greater than 90 days. This decomposition turns out to be 445 considerably different for the CTRL and UBR0 simulations. The differences illustrate 446 the effect that the mean winds have on wave propagation. Since the tropical waves 447 responsible for driving the RMMC have phase speeds that are comparable to the wind 448 speeds in the tropics and subtropics, it is not surprising that the EP-flux divergence 449 pattern is sensitive to the mean winds.

The EP-flux and divergence for the DJF07 season, decomposed into stationary, westward, and eastward propagating waves are shown in Figure 9. The CTRL simulation is on the left and the UBR0 simulation is on the right. As we have seen, the wave heating source was confined primarily in the southern hemisphere (Fig. 1), the mean winds (Fig. b) have magnitude $\sim 5 \text{ ms}^{-1}$ or less in the tropics, (westward in the lower troposphere, eastward in the upper troposphere), and the subtropical jet is much stronger in the

northern hemisphere (30 ms⁻¹) than the southern hemisphere (15 ms⁻¹). Thus we see, for 456 457 stationary and westward waves in both simulations, the EP flux emanating upward from 458 the southern tropics and subtropics, with some propagation toward and across the 459 equator. For the CTRL simulation the strong northern subtropical jet appears to allow 460 eastward wave propagation out of the northern hemisphere as well. Relative to the UBR0 461 simulation, the CTRL simulation has increased the EP flux divergence for stationary and 462 eastward waves and decreased the divergence for westward waves. The decrease in 463 westward wave flux is likely a manifestation of the Charney-Drazin criterion (Charney 464 and Drazin, 1961) - the subtropical jets have mean wind speed that is too large for 465 propagation of the Rossby waves that are excited. The nonzero mean winds also allow 466 for Rossby waves that propagate eastward relative to the ground but westward relative to the mean flow. 467

The eastward wave spectrum also contains Kelvin waves. The EP-flux of Kelvin waves point downward over the equator (since by (2) F_z is upward flux of westward momentum) and, when dissipating, will produce a positive EP-flux divergence. There is a similar region of positive EP-flux divergence and downward flux for eastward waves in both simulations. The Kelvin waves present in these simulations have larger phase speeds than the Rossby waves, and are less affected by the relatively small mean wind speeds in the tropical troposphere.

475 The RMMC that is driven by the EP-flux divergences for the various wave 476 components from the CTRL and UBR0 simulations are shown in Figure 10. The relative 477 strengths of the residual mean vertical velocity are essentially the same as the relative 478 strengths of the EP-flux divergences shown in Fig. 9. The upwelling driven by stationary 479 waves (Fig 10, top) is stronger in the CTRL than in the UBRO simulation. The EP-flux 480 divergence pattern for CTRL is shifted poleward from the divergence pattern for UBR0. 481 Thus, the upward residual mean velocity is shifted with a maximum that lies south of the 482 equator. Conversely, the upwelling driven by westward waves (Fig 10, middle) is weaker

in CTRL than in UBR0 due to the suppression of westward Rossby wave propagation. 483 484 The Kelvin waves (Fig. 10, bottom) force downward residual mean velocity over the 485 equator. This is most evident in UBR0, whereas eastward Rossby waves are also present 486 in CTRL. The upwelling driven by the eastward Rossby waves largely cancels the down-487 welling driven by the Kelvin waves in this case. Since the downward residual mean 488 velocity driven by the Kelvin waves is more narrowly confined to the equator than the 489 upward velocity driven by the Rossby waves, the total RMMC obtained from the 490 superposition of wave components may actually show a local minimum over the equator. 491 Such a local minimum is evident in the residual mean vertical velocities shown for DJF07 492 in Fig. 4.

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494 e) Interannual variability

The QBO may influence the RMMC by affecting the waveguide for the tropical waves. The CTRL simulation includes both phases of the QBO. The mean winds shown in Fig. 5 have a QBO westerly phase overlying the TTL from August 2006 through June 2007. The question is whether the QBO signal penetrates deep enough to influence the TTL wave propagation in a significant manner.

There may be interannual variability in the TRMM heating that is used to force waves in our model. To determine if the interannual variability in the RMMC of the CTRL simulation is directly related to wave propagation through the QBO winds, we performed an additional experiment. The model was run using winds from 2006 and TRMM heating from 2007, and conversely, using winds from 2007 and TRMM heating from 2006. Any difference between these runs and the CTRL simulation can only be attributed to a change in the background winds and temperatures.

507 Examination of the four cross sections of EP-flux and RMMC in each season for 508 each of the cases of 2006/2007 TRMM heating paired with 2006/2007 winds (not shown here) did not demonstrate convincing differences amongst them that can definitely be attributed to differences in QBO tropical winds. This is not surprising – if one covers the top panel of Fig. 5 with their hand to block out winds above the TTL, it is difficult to discern a QBO pattern in the winds below. Based on the EP-fluxes and divergences shown in previous figures, we note that Rossby waves do not propagate much higher than the TTL.

515 However, Kelvin waves can readily propagate into the QBO easterly phase, so it 516 is possible that there is a difference in the Kelvin wave signature in the RMMC. From 517 what was discussed above, the Kelvin waves drive a downward circulation over the 518 tropics, potentially forming a local minimum in the RMMC over the equator. There does 519 seem to be a QBO signal in such a Kelvin wave contribution to the RMMC pattern. 520 Figure 11 shows the vertical residual mean velocity averaged over the TTL for the four 521 separate pairings of mean winds and heating in our model runs. The two solid curves in 522 the figure are obtained using observed winds from SON 2006 when QBO westerly winds 523 are in the lower stratosphere. These two curves, produced using TRMM heating from 524 different years, have a noticeable local minimum over the equator, indicating that a 525 circulation driven by dissipating Kelvin waves is present. The dashed curves, obtained 526 using observed winds from 2007 show a depression over the tropics of much smaller 527 magnitude. A plausible explanation for this is that the westerlies present in 2006 inhibit 528 some Kelvin waves from propagating higher into the stratosphere, causing them to 529 dissipate just above the TTL and producing a significant contribution to positive EP-flux 530 divergence.

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532 f) Contributions to upwelling from synoptic waves

Finally, we compare the magnitude of the upwelling in the TTL driven by tropicalwaves to the upwelling driven by synoptic waves. Synoptic waves in our model grow off

535 an unstable configuration of the zonal mean subtropical jet, but the strength of these 536 waves depend on how strongly we clamp the model background winds to the observed 537 zonal means. All studies using our model showed that synoptic waves make a small 538 contribution to upwelling through the TTL. However, our model does not incorporate 539 most of the physical processes in the troposphere, so it is difficult to verify the validity of 540 our simulations of synoptic wave dynamics. Here instead the MERRA 3-hourly 541 assimilated product for 2006-2007 is analyzed following the approach already described. 542 Namely, we calculated the EP-flux divergence spectrum for each month, which was then 543 used to force the TEM model to calculate the RMMC response. To isolate the synoptic 544 waves as much as possible, we used only the transient portion of the spectrum - all waves 545 with period less than 30 days. We assume that most of these waves are synoptic, although the results of Section 3d above show that a portion of the transient wave 546 547 spectrum arises from tropical waves forced by latent heating. Even without being able to 548 make a clean separation between synoptic and tropical waves using MERRA analysis, we 549 obtain the same conclusion as with our own model simulations, that synoptic waves make 550 a negligible contribution to TTL upwelling.

551 The monthly average of the upwelling calculated using the TEM model forced 552 with the EP-fux divergence from MERRA transient waves is shown in the top panel of 553 Figure 12. Compared to the residual vertical mean velocities shown in the upper left 554 panel of Fig. 8, the upwelling forced by MERRA transient waves is weaker than what is 555 forced by tropical waves by a factor of two or more. Also, the altitude where the peak 556 magnitude occurs is shifted lower into the troposphere. Upwelling in the TTL forced by 557 synoptic waves is estimated to be less than .2 mm/s. We see that synoptic waves 558 primarily contribute to upwelling into the lower part of the TTL whereas tropical waves 559 are far more effective at producing upwelling all the way through the TTL. It is also 560 interesting to note that the upwelling forced by the MERRA transient waves has a semi-561 annual variation that peaks near equinox and has a minimum during solstice.

The EP-flux divergence for MERRA transient waves and the RMMC obtained from the TEM model for SON 2007 are shown in the middle and bottom panels of Fig. 12, respectively. Compared to Figs. 3 and 4, the peak in the transient wave EP-flux convergence has similar magnitude, but occurs 5 km lower than and around 5° poleward of the tropical wave EP-flux convergence. It is for these reasons that the upwelling forced by MERRA synoptic waves is weaker and lower than the upwelling forced by tropical waves.

569

570 **4.** Conclusions

571 We have used a model forced with heating derived from TRMM rainfall rate 572 observations to demonstrate that equatorial waves forced by latent heating drive most of 573 the observed residual mean upwelling across the TTL. Our model simulations produce a 574 residual mean vertical velocity that is essentially the same as values derived from 575 observations within 15° of the equator. Our experimental design has enabled us to isolate 576 equatorial waves from extratropical synoptic waves and determine their relative 577 contribution to driving the upward vertical component of the RMMC. Our estimates 578 show that the upwelling in the TTL driven by equatorial waves is much stronger than the 579 upwelling driven by synoptic waves. The synoptic wave-driven circulation estimated 580 from MERRA reanalysis is confined to much lower altitudes than the tropical wave-581 driven circulation in our simulations because tropical waves can propagate higher into the 582 TTL and tropical lowermost stratosphere. The EP-flux convergence from synoptic scale 583 waves resides mainly in the troposphere and is therefore not positioned to force residual 584 mean upwelling across the TTL.

Recent work has highlighted the role of synoptic waves in driving future trends in the RMMC in climate models (Shepherd and McLandress, 2011), and the mechanism involves a rising critical level for these waves due to greenhouse-gas-related temperature

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changes. Our work does not address these future changes, but examines conditions in the recent past. However, we note that small changes in tropical tropospheric winds on the order of only 5 m s⁻¹ have large effects on the tropical wave propagation and TTL upwelling in our model study. This suggests the possibility that small wind biases or differences among climate models could lead to differing conclusions on attribution of wave driving of the TTL upwelling in different models.

594 We note the TTL upwelling forced by tropical waves in our simulations is largely 595 confined to latitudes within 20° of the equator. Clearly, other waves with extratropical 596 origins and EP flux divergence at subtropical and extratropical latitudes are also 597 important for forcing the full equator-to-pole Brewer-Dobson circulation throughout the 598 stratosphere that is inferred from tracer observations. We highlight here the importance 599 of tropical waves in driving the upwelling within and through the TTL at altitudes, above 600 the influence of latent heating due to convection and below the 70hPa reference level that 601 has been used in many previous model studies (e.g. Butchart et al., 2010). Our results 602 (Figure 8) show that the tropical upwelling driven by tropical quasi-stationary Rossby 603 waves is weak above 100 hPa.

604 We have also examined the variability of the TTL upwelling throughout a two 605 year simulation that includes two phases of the QBO. Although the TRMM heating that 606 drives the waves has considerable seasonal variability, we found that the circulation 607 driven by the wave response to this heating in a run where the background zonal mean 608 zonal winds were held fixed to zero did not have a regular annual cycle. A strong annual 609 variation in the upwelling is produced when the waves propagate through seasonally 610 variable winds. The maximum upwelling value, averaged within 15° degrees of the equator, varies between .4 mm s^{-1} and .8 mm s^{-1} . When the upwelling is additionally 611 averaged between 90-190 hPa, it varies between .2 mm s⁻¹ and .4 mm s⁻¹. The seasonal 612 613 variation in the subtropical jets modulates the hemisphere in which the strongest EP-flux 614 divergence occurs. A period between June and October when weak mean easterlies occur over the equator corresponds to weaker wave propagation, EP-flux divergence and RMMC. The variability in TTL upwelling appears to result from relatively large changes in Rossby wave propagation due to only weak variation in tropical wind speed of less than 10 m s⁻¹, with weak westward winds of around 5 m s⁻¹ corresponding to the period of weakest upwelling. In addition, variability in the heating appears to be responsible for producing anomalies in the interannual variability of the upwelling (for example, the large upwelling in January 2007).

622 The mean winds also modify the frequency spectrum of the tropical waves. When 623 the background winds are zero, the portions of the wave spectrum responsible for driving 624 the RMMC are equally split between westward waves with low frequency and stationary 625 waves. With observed background winds, stationary waves are primarily responsible for 626 driving the RMMC. There is still significant wave driving contribution from the transient 627 waves with low frequency. This contribution seems equally split or slightly weighted 628 toward eastward traveling Rossby waves. A downward circulation component due to 629 Kelvin wave driving is detectible in our experiments. Our limited two-year run suggests 630 that the QBO in the lower stratosphere may modulate the Kelvin wave driving near the 631 tropopause, and that this may be the only way that the QBO winds directly affect the TTL 632 upwelling due to tropical waves. Figure 13 summarizes these findings, illustrating the 633 annual variability of the upwelling in relation to the annual variability of the mean winds, 634 and the interannual influence of the QBO on the circulation associated with Kelvin wave 635 damping. We note that while our simplified model approach allows us to isolate the 636 effects of tropical and extratropical waves on TTL upwelling, it cannot address the causes 637 of the annual cycle in the tropical zonal wind itself, which in the real atmosphere results 638 from a balance of eastward and westward accelerations due to tropical waves and the 639 Hadley cell overturning (e.g. Kraucunas and Hartman, 2005).

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29

754 **Figure Captions.**

30

Figure 1. TRMM heating at 7 km and over four seasons in 2007.

Figure 2. Top: The eddy component of the temperature field and zonal and vertical wind components averaged from 10°S-10°N and 2007. The TTL is bounded by the horizontal red lines. Vertical scale of the vectors is exaggerated by a factor of 1000 in this and the following figures. Middle: The eddy component of the temperature field at 26 hPa averaged over 2007. Bottom: The temperature field at 100 hPa averaged over 2007.

Figure 3. EP-flux (vectors) and its divergence (contours) for equatorial waves forced by heating derived from TRMM rainfall, averaged over four seasons in 2007. The results for the simulation with zero background winds are in the left column and for observed background winds in the right column.

Figure 4. Residual mean circulation (vectors) and the residual mean vertical velocity (contours) for equatorial waves forced by heating derived from TRMM rainfall, averaged over four seasons in 2007. The results for the simulation with zero background winds are in the left column and for observed background winds in the right column.

Figure 5. Zonal and monthly mean zonal winds from NCEP for 2006-2007. Top: Vertical profiles averaged from 15°S-15°N, (contour interval 3 ms⁻¹); Bottom: Latitude cross sections averaged from 90-190 hPa (contour interval 4 ms⁻¹).

Figure 6. Zonal and monthly mean vertical component of EP-flux for 2006-2007. Top
row: Vertical profiles averaged from 15°S-15°N; Bottom row: Latitude cross sections
averaged from 90-190 hPa. Left column: CTRL simulation; Right column: UBR0
simulation.

- **Figure 7.** Same as Fig.6, for EP-flux divergence.
- **Figure 8.** Same as Fig. 6, for vertical residual mean velocity.

Figure 9. EP-flux (vectors) and its divergence (contours) for the DJF 2007 CTRL
simulation (left column) and UBR0 simulation (right column). Top row: Stationary
waves only. Middle row: Westward waves only. Bottom row: Eastward waves only.

Figure 10. Same as Fig. 9, for the residual mean vertical velocity.

Figure 11. Residual mean vertical velocity averaged over SON and the TTL altitudes.
Black curves are for TRMM heating from 2006, red from 2007. Solid curves are for
observed background winds from 2006, dashed from 2007. The solid curves show a local
minimum near the equator that likely is the result of Kelvin wave forcing.

Figure 12. Response of the TEM model to MERRA transient wave forcing. Top:
Residual mean vertical velocity, averaged from 15°S-15°N; Middle: EP-flux (vectors) and
its divergence (contours); Bottom: Residual mean circulation (vectors) and residual
mean vertical velocity (contours).

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790 Figure 13. Schematic illustrating the annual and quasi biennial influence of tropical 791 waves and TTL upwelling. The color background shows 15S-15N zonal mean winds 792 versus time and height over the two years 2006-2007, with pink shades westerly and blue 793 shades easterly. The TTL region is denoted with dashed black lines. Purple arrows 794 indicate the stationary equatorial Rossby wave-driven upwelling, which is stronger 795 during boreal winter when upper troposphere winds are more westerly, and weaker in 796 boreal summer when winds are more easterly. The smaller red arrows indicate Kelvin 797 wave-driven downwelling, which generally occurs higher in the stratosphere, but can 798 influence the TTL when the QBO westerly winds reach the lower stratosphere.

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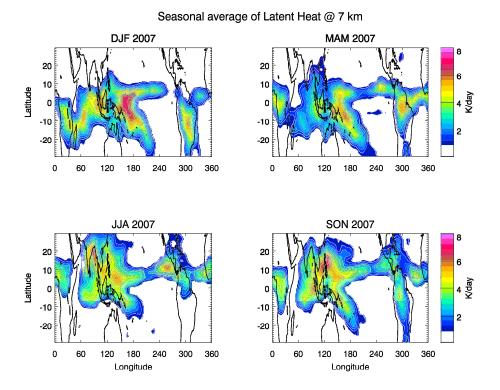


Figure 1. TRMM heating at 7 km and over four seasons in 2007.

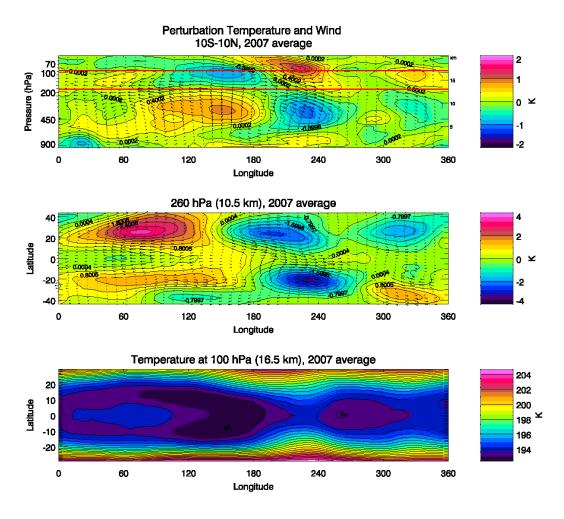
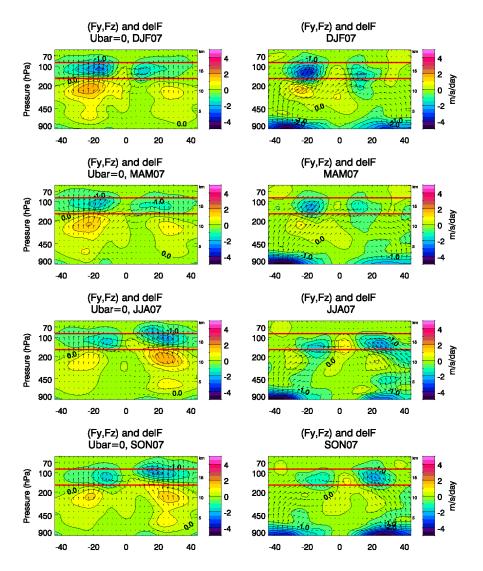


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811 Figure 3. EP-flux (vectors) and its divergence (contours) for equatorial waves forced by 812 heating derived from TRMM rainfall, averaged over four seasons in 2007. The results for 813 the simulation with zero background winds are in the left column and for observed 814 background winds in the right column.

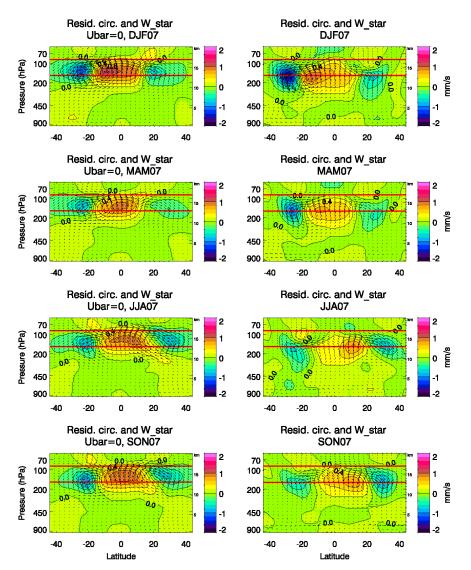


Figure 4. Residual mean circulation (vectors) and the residual mean vertical velocity
(contours) for equatorial waves forced by heating derived from TRMM rainfall, averaged
over four seasons in 2007. The results for the simulation with zero background winds are
in the left column and for observed background winds in the right column.

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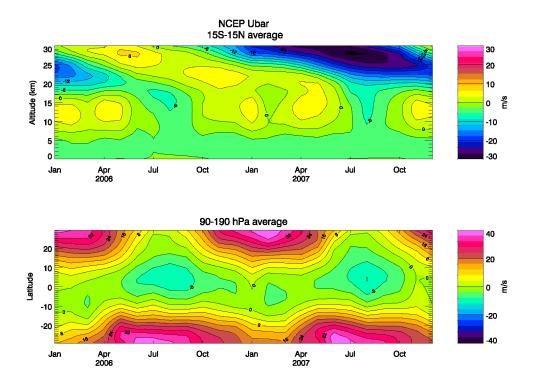


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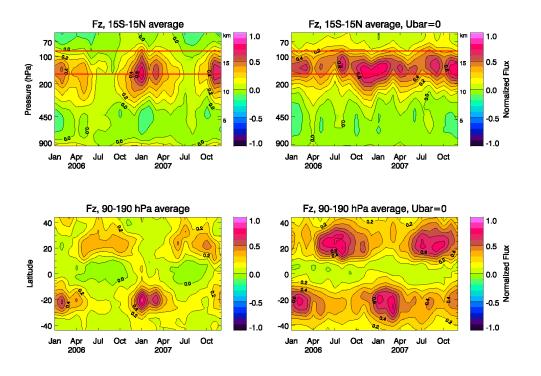
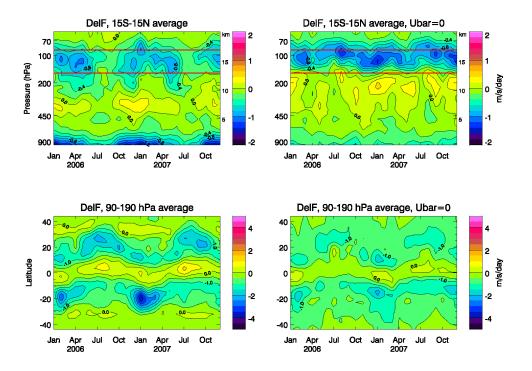
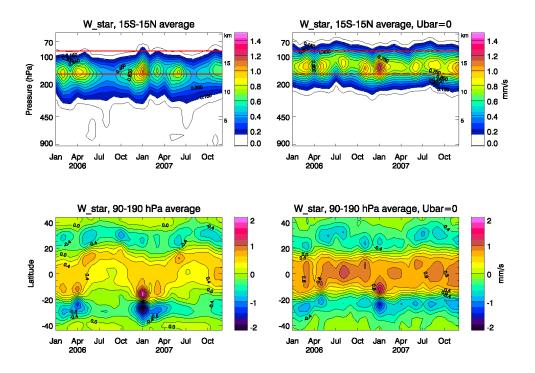


Figure 6. Zonal and monthly mean vertical component of EP-flux for 2006-2007. Top
row: Vertical profiles averaged from 15°S-15°N; Bottom row: Latitude cross sections
averaged from 90-190 hPa. Left column: CTRL simulation; Right column: UBR0
simulation.



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833 **Figure 7.** Same as Fig.6, for EP-flux divergence.



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836 **Figure 8.** Same as Fig. 6, for vertical residual mean velocity.

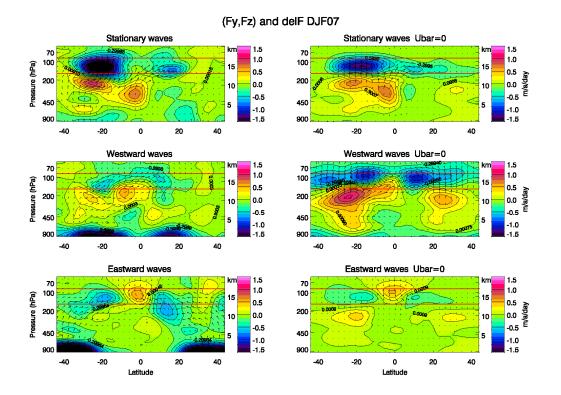
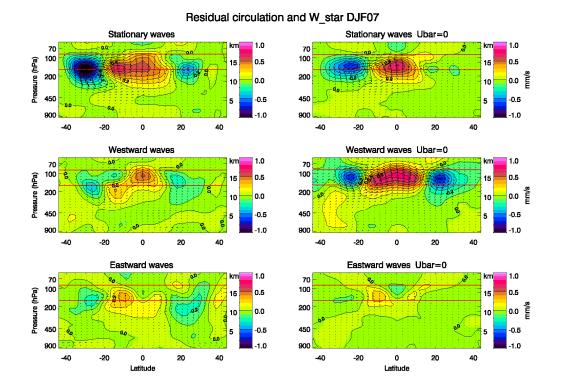


Figure 9. EP-flux (vectors) and its divergence (contours) for the DJF 2007 CTRL
simulation (left column) and UBR0 simulation (right column). Top row: Stationary
waves only. Middle row: Westward waves only. Bottom row: Eastward waves only.

842



845 **Figure 10.** Same as Fig. 9, for the residual mean vertical velocity.

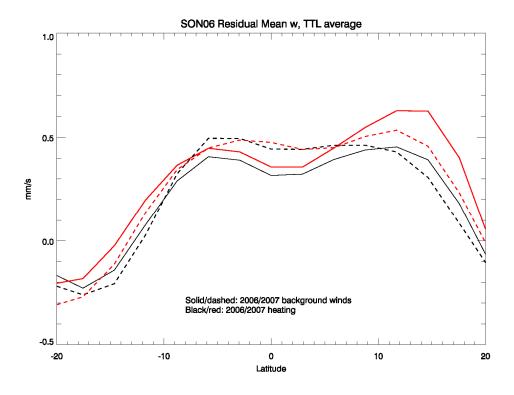
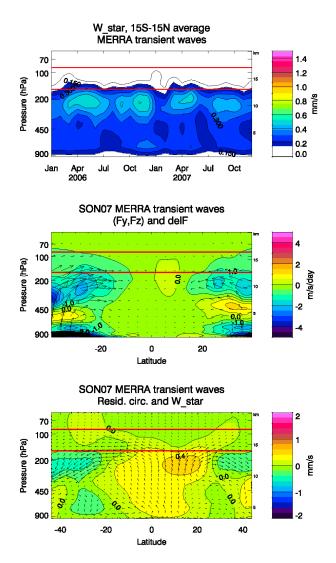
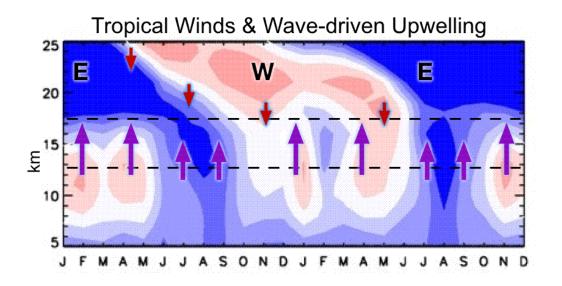


Figure 11. Residual mean vertical velocity averaged over SON and the TTL altitudes.
Black curves are for TRMM heating from 2006, red from 2007. Solid curves are for
observed background winds from 2006, dashed from 2007. The solid curves show a local
minimum near the equator that likely is the result of Kelvin wave forcing.



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Figure 12. Response of the TEM model to MERRA transient wave forcing. Top:
Residual mean vertical velocity, averaged from 15°S-15°N; Middle: EP-flux (vectors) and
its divergence (contours); Bottom: Residual mean circulation (vectors) and residual
mean vertical velocity (contours).



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860 Figure 13. Schematic illustrating the annual and quasi biennial influence of tropical 861 waves and TTL upwelling. The color background shows 15S-15N zonal mean winds versus time and height over the two years 2006-2007, with pink shades westerly and blue 862 863 shades easterly. The TTL region is denoted with dashed black lines. Purple arrows 864 indicate the stationary equatorial Rossby wave-driven upwelling, which is stronger during boreal winter when upper troposphere winds are more westerly, and weaker in 865 866 boreal summer when winds are more easterly. The smaller red arrows indicate Kelvin 867 wave-driven downwelling, which generally occurs higher in the stratosphere, but can 868 influence the TTL when the QBO westerly winds reach the lower stratosphere.