

1 **Unexpected DE3 tide in the southern summer**
2 **mesosphere**

Erich Becker¹, Jens Oberheide²

Corresponding author: E. Becker, Northwest Research Associates, 3380 Mitchell Ln, Boulder, CO 80301, USA (erich.becker@nwra.com)

¹NorthWest Research Associates, Boulder, CO, USA

²Department of Physics and Astronomy, Clemson University, Clemson, SC, USA

Key Points.

- SABER satellite observations and MERRA-2 reanalysis show a significant DE3 tidal component in the southern summer mesosphere.
- The HIAMCM with resolved gravity waves confirms these observations and shows that the DE3 extends up to ~ 90 km at high latitudes.
- The summer mesospheric DE3 gives a 10-20% contribution to the eastward EPF divergence that drives the equatorward residual circulation.

3 Simulation of the January 2017 period us-
4 ing a gravity-wave resolving global circulation
5 model (HIAMCM) reveals a predominant east-
6 ward propagating diurnal tide with zonal wavenum-
7 ber three (DE3) in the southern summer meso-
8 sphere from about 60 to 90 km height at mid-
9 dle to high latitudes. We provide observational
10 evidence based on MERRA-2 reanalysis and
11 SABER satellite observations for the validity
12 of this result. The attenuation of the DE3 be-
13 low the mesopause generates a significant east-
14 ward Eliassen-Palm flux divergence that con-

15 tributes to the residual circulation. We also
16 show that the diurnal tide in the northern sum-
17 mer mesosphere likely consists of mainly east-
18 ward propagating components. These findings
19 contradict the common perception of a weak
20 diurnal tide in the summer mesosphere.

1. Introduction

21 Thermal tides have long been known to represent the strongest wave-related wind and
22 temperature perturbations in the mesosphere and lower thermosphere (MLT) [e.g. *Forbes*,
23 1984; *Akmaev*, 2001]. The tides relevant in the MLT are generated by the daily cycle of
24 UV absorption by stratospheric ozone, as well as by the absorption of infrared solar
25 radiation by tropospheric water vapor and clouds. This gives rise to so-called migrating
26 tides, that is, tidal components that propagate synchronously with the sun from East
27 to West, such as the diurnal tide with zonal wavenumber $s = 1$ (DW1) and the semi-
28 diurnal tide with $s = 2$ (SW2). Additional tidal forcing is due to the daily cycle of
29 deep moist convection in the intertropical convergence zone. Since this tidal forcing is
30 localized in certain geographical regions, so-called non-migrating tides develop that do not
31 propagate synchronously with the sun [*Hagan and Forbes*, 2002; *Zhang et al.*, 2010a, b].
32 Non-migrating tides are also formed by the interactions of migrating tides with planetary
33 waves and the mean flow [*Lieberman et al.*, 2014; *Achatz et al.*, 2008]. The most prominent
34 non-migrating component is the diurnal, eastward propagating tide that has $s = 3$, the
35 so-called DE3.

36 Current understanding of tides is based on linear theory [*Chapman and Lindzen*, 1969],
37 and on linear numerical models that include realistic forcing and background atmospheres
38 [*Hagan and Forbes*, 2002; *Achatz et al.*, 2008; *Oberheide et al.*, 2009]. Such linear models
39 can describe many observations of tides in the MLT, and they agree with the tides simu-
40 lated by general circulation models (GCMs) with parameterized gravity waves (GWs) [e.g.,
41 *Smith*, 2012; *Ward et al.*, 2010; *Liu et al.*, 2010; *Pedatella et al.*, 2016; *Vitharana et al.*,

2019]. The morphology of tides can be summarized as follows. The DW1 is the most prominent component at low and subtropical latitudes up to about 90 km. In the 90-130 km altitude range at these latitudes, the DE3 has the largest amplitude in comparison to all other tidal components. The tides have smaller amplitudes at middle and high latitudes. According to ground-based measurements, the diurnal tides are most prominent up to about 85-90 km, while semi-diurnal tides account for the strongest tidal variations in the mesopause region [e.g. *Lübken et al.*, 2011; *Kishore Kumar et al.*, 2014]. According to linear models and conventional GCMs, the DW1 should be very weak in the mesosphere at middle to high latitudes, while the SW2 becomes significant around and above the mesopause at these latitudes [e.g. *Smith*, 2012, Figs. 9 and 13].

In the present study we investigate the question whether the observed diurnal variations in the extratropical summer mesosphere below ~ 85 -90 km are possibly caused by tidal components other than the DW1. For this purpose we use a GCM with explicit simulation of GWs, as well as reanalysis and satellite observations. In Sec. 2 we describe the model and define our tidal analysis. In Sec. 3 we compare tidal variations from the model and reanalysis, and we estimate the relevance of the DE3 for the residual circulation. Section 4 presents a new analysis of SABER temperatures. Our conclusions are presented in Sec. 5.

2. Model and tidal analysis

We employ the High Altitude Mechanistic general Circulation Model (HIAMCM). This model is based on a standard spectral dynamical core that is extended by non-hydrostatic dynamics and thermodynamics for variable composition. It is run at a T256 spectral hor-

63 izontal resolution and with 280 atmospheric layers extending up to 4×10^{-9} hPa ($z \sim 400$ -
64 500 km). The HIAMCM includes radiative transfer, water vapor transport, latent heating,
65 full topography, a simple slab ocean model, the full surface energy budget, and simple
66 representations of ion drag in the thermosphere. Macro-turbulent vertical and horizon-
67 tal diffusion is represented by the Smagorinsky scheme, with both diffusion coefficients
68 depending on the Richardson number. This diffusion scheme accommodates molecular
69 viscosity and heat conduction.

70 The HIAMCM can be nudged to the three-hourly Modern-Era Retrospective analysis
71 for Research and Applications version 2 (MERRA-2) [*Bosilovich et al.*, 2015; *Gelaro et al.*,
72 2017]. This nudging is performed in spectral space and is restricted to the large-scale flow
73 such that the resolved GWs are not directly affected by the nudging. More specifically,
74 we interpolate the MERRA-2 wind and temperature fields to the terrain-following grid of
75 the HIAMCM and compute the MERRA-2 spectral representations of relative vorticity,
76 horizontal divergence, and temperature. This allows nudging the HIAMCM in spectral
77 space. Moreover, all postprocessing routines developed for the HIAMCM can be applied
78 to MERRA-2 as well. The HIAMCM does not include parameterization of GWs and has
79 an effective resolution of a horizontal wavelength of ~ 200 km. Detailed information about
80 the HIAMCM can be found in *Becker and Vadas* [2020], *Becker et al.* [2022a], and *Becker*
81 *et al.* [2022b].

82 In the following we analyze time series for 1-20 January 2017 and 6-20 July 2006. Before
83 applying the usual tidal decomposition, we first compute average daily cycles in spectral
84 space. An average daily cycle from the HIAMCM is defined as follows. We compute

85 temporal averages of the universal time intervals 23:30-00:30, 01:00-02:00, 02:30-03:30, ...,
86 and 22:00-23:00, taking all model days of the respective period into account. This leads
87 to a times series with 16 time stamps centered at universal times 00:00, 01:30, 03:00, ...,
88 and 22:30. Given that the HIAMCM spectral coefficients are saved every 10 minutes, each
89 time stamp of the average daily cycle for 1-20 January 2017 represents an average over
90 $7 \times 20 = 140$ snapshots. This number is $7 \times 15 = 115$ for 6-20 July 2006. In the case of
91 MERRA-2, snapshots are available every 3 hours. Hence, average daily cycles have 8 time
92 stamps at universal times of 00:00, 03:00, ..., and 21:00 UT. Each time stamps represents
93 an average over 20 (15) snapshots for the January 2017 (July 2006) period. The total tide
94 is defined as the average daily cycle minus its 24 h average.

3. Results from the model and reanalysis

3.1. Tidal structure and amplitudes

95 Figure 1 illustrates the total temperature tide at 55°S for 1-20 January 2017 from the
96 HIAMCM (left column) and MERRA-2 (right column). The first row shows longitude-
97 time plots at 0.04 hPa ($z \sim 75$ km), while the second and third rows show longitude-
98 height plots at 00:00 UT and 12:00 UT, respectively. The upper level of MERRA-2 is
99 indicated by horizontal black lines in panel c-f to facilitate the comparison between the
100 left and right panels. Figures 1a and b indicate a significant DE3 in the southern summer
101 mesosphere. This DE3 is superposed with other components, particularly a DW1. The
102 DE3 is more prominent in the HIAMCM, while the DW1 is more prominent in MERRA-2.
103 The longitude-height plots from the model reveal that the DE3 extends from about 0.1
104 (60 km) to 0.001 hPa (90 km). The tidal structure in MERRA-2 (panel d,f) agrees with

105 that from the model below 0.015 hPa. In particular, a predominant DW1 from about 5 to
106 0.3 hPa is seen in both data sets with similar amplitudes and phases. Overall, the total
107 tide is more structured in the HIAMCM.

108 Figure 2 shows temperature amplitudes of individual tidal components for 1-20 January
109 2017. The HIAMCM (MERRA-2) results are shown in the left (right) column. Panels
110 a and b show similar DW1 amplitudes below 0.015 hPa, except for the mesosphere at
111 middle and high latitudes where the DW1 has larger amplitudes in MERRA-2. The
112 DW1 furthermore exhibits maxima in the MLT over the equator and around 30° to 40°
113 latitude in either hemisphere. This behavior is well known from other studies [e.g. *Smith*,
114 2012, her Fig. 8]. The SW2 from the HIAMCM (panel c) exhibits subtropical maxima in
115 the lower thermosphere, but is also significant at middle to high latitudes in the upper
116 mesosphere. MERRA-2 shows larger SW2 amplitudes in the northern lower mesosphere
117 than the HIAMCM. The SW2 amplitudes in Fig. 2c,d are similar in the stratopause region
118 at low latitudes.

119 The third row of Fig. 2 shows the amplitudes of the eastward propagating, non-migrating
120 tidal components. Colours show the sum of DE1, DE2, and DE3, while white contours
121 show the DE3. Both the HIAMCM and MERRA-2 indicate a tropical maximum of the
122 DE3 below 0.015 hPa. In the HIAMCM (panel e), this maximum transitions into a broad
123 maximum in the lower thermosphere that extends into the subtropics, which is well known
124 from analysis of satellite observations [e.g., *Kumari et al.*, 2020] and GCMs [e.g., *Smith*,
125 2012, Fig. 14]. Even though the DE3 gives the main contribution in this regime, the DE1
126 and DE2 components are also significant. Around 0.0001 hPa, the combined amplitude

127 of the DE components at tropical and subtropical latitudes is significantly larger than the
128 DW1 amplitude.

129 The main finding of this study is that DE tides exhibit a pronounced maximum in
130 the southern summer upper mesosphere at middle to high latitudes, with the DE3 giving
131 the predominant contribution (Fig. 2e). The HIAMCM and MERRA-2 both show that
132 the DE3 is significant in the mesosphere below 0.015 hPa from about 20°S to 60°S. The
133 HIAMCM indicates that this maximum shifts toward the pole with increasing altitude
134 and has a maximum at 60°S and 0.003 hPa (about 85 km).

135 This result is quite surprising given the fact that the DE3 is usually found only at and
136 above the mesopause at low latitudes. More specifically, the DE3 is considered to be the
137 superposition of a Kelvin wave-like broad symmetric mode that maximizes above 100 km
138 over the equator and an anti-symmetric tidal mode that maximizes around $\pm 20^\circ$ latitude
139 and 95 km [*Oberheide and Forbes, 2008*], with both modes exchanging energy in the
140 stratosphere/mesosphere when propagating upward [*Zhang et al., 2012*]. The 20°S/N
141 amplitude maxima in Fig. 2e in the mesosphere with the transition into a broad amplitude
142 maximum symmetric about the equator at higher altitudes is thus what is expected from
143 tidal theory and observations. However, the presence of the DE3 at 60°S and 0.003
144 hPa is unexpected and cannot be explained through higher-order Hough modes. This
145 is because the second symmetric and antisymmetric modes both peak equatorward of
146 30° latitude, and the vertical wavelengths of the third symmetric and antisymmetric modes
147 are well below 10 km and as such too small for these modes to propagate upward from
148 the troposphere.

149 *Lübken et al.* [2011] analyzed lidar temperature measurements performed at the station
150 of Davis (69°S, 78°E, Antarctica) during January 2011. They found a significant diurnal
151 temperature tide in the upper mesosphere with an amplitude of at least 6 K at ~ 85 km
152 (see Fig.2 in their paper). They mentioned that conventional models show much weaker
153 tidal amplitudes in this region. A DE3 maximum of about 6 K at 69°S and 0.003 hPa
154 (~ 85 km) as simulated by the HIAMCM (Fig. 2e) is quantitatively consistent with the
155 lidar result. Moreover, when considering Figs. 1c and e, also the phase of this diurnal
156 variation with maximum temperatures around local noon at ~ 85 –90 km agrees with the
157 lidar result, even though Fig. 1 shows results for 55°S.

158 The first row in Fig. 3 illustrates the temperature tide at 55°N for 6-20 July 2006
159 from the HIAMCM and MERRA-2. Comparison of Fig. 3a,b to Fig. 1a,b indicates that
160 eastward propagating tidal components are less prominent in the northern than in the
161 southern summer mesosphere. As a result, the DW1 is of stronger relative importance in
162 both the HIAMCM and MERRA-2. Figure 3c,d show DW1 and SW2 tidal temperature
163 amplitudes for 6-20 July 2006 from the HIAMCM. These amplitudes are similar to that
164 for January when comparing the respective winter and summer hemispheres. In partic-
165 ular, the DW1 is small in the northern summer mesosphere at middle to high latitude.
166 Figure 3e,f show the DE amplitudes from the HIAMCM and MERRA-2. Strong DE am-
167 plitudes are seen in the northern summer mesosphere. However, these components are less
168 significant than during January.

169 The HIAMCM shows a maximum north of 60°N between 0.01 and 0.003 hPa of about
170 3 K due to the sum of DE1, DE2, and DE3, where the DE3 gives a contribution of at

171 most 1 K (Fig. 3e). This result agrees with lidar observations of *Gerding et al.* [2013]
172 during June and July from 2010 to 2013 at the station of Kühlungsborn (54°N, 11°E).
173 These authors found maximum diurnal variations at ~ 85 km of a few K (see Fig. 4 in
174 their paper), which is weaker than the aforementioned result for Antarctica. According
175 to Fig. 3, the DE components can explain these diurnal tidal variations.

176 We note that the DE components account for the main diurnal variations in the southern
177 winter stratopause region from about 1 to 0.1 hPa at middle to high latitudes (Figs. 3c-
178 f). This feature was also found by *Sakazaki et al.* [2012] in both satellite observations
179 and reanalyses. We speculate that these authors did not discover DE components in the
180 summer mesosphere because their analysis was restricted to altitudes below ~ 65 km.

3.2. Relevance for the general circulation

181 Figures 4a-d show the zonal-mean circulation from the upper stratosphere to the lower
182 thermosphere from the HIAMCM for 1-20 January 2017 (left column) and 6-20 July 2006
183 (right column). The HIAMCM simulates reasonably realistic temperatures and zonal
184 winds (Fig. 4a,b). This includes the cold summer mesopause and the transition from
185 westward to eastward flow above the temperature minimum, the subtropical mesospheric
186 jet in the winter hemisphere, as well as eastward winds at high latitudes in the win-
187 ter MLT. There are important hemispheric differences when comparing July to January.
188 These include a stronger eastward flow and stronger westward Eliassen-Palm flux (EPF)
189 divergence in the winter mesosphere, stronger absolute EPF divergence in the upper meso-
190 sphere and a stronger summer-to-winter pole residual circulation (Fig. 4c,d), and a colder
191 summer polar mesopause (Fig. 4a,b). These hemispheric differences are consistent with

192 satellite observations and the interhemispheric coupling mechanism [e.g. *Körnich and*
193 *Becker*, 2010; *Smith*, 2012; *Karlsson and Becker*, 2016]. There is stronger eastward EPF
194 divergence in the winter mesopause region during July than during January. According
195 to previous studies [e.g., *Becker and Vadas*, 2018; *Vadas and Becker*, 2019; *Harvey et al.*,
196 2022; *Becker et al.*, 2022b], this hemispheric difference is caused by stronger secondary
197 GWs in the winter MLT for a stronger polar vortex. Also the westward EPF divergence in
198 the summer lower thermosphere is stronger during July. As a result of these hemispheric
199 differences, the reversed residual circulation cell in the lower thermosphere [*Smith et al.*,
200 2011] is stronger during July and extends from pole to pole.

201 Figures 4e,f show the EPF divergence due to the resolved GWs in the summer MLT
202 (colors). We compute the GW EPF divergence by subtracting the EPF divergence that
203 is due to planetary and synoptic scales. The latter is defined by applying a triangular
204 spectral truncation at a wavenumber of 30 to the model output. The so-defined GW
205 EPF divergence exceeds $120 \text{ ms}^{-1}\text{d}^{-1}$ at 50°N to 60°N around 0.001 hPa during July,
206 which is comparable to estimates from GW schemes [e.g. *Fomichev et al.*, 2002, their Fig.
207 10]. In the HIAMCM, however, this GW drag is too high in altitude by $\sim 5 \text{ km}$. As a
208 result, also the summer mesopause and the zonal wind reversal are too high in altitude by
209 $\sim 5 \text{ km}$. Superposed in Figs. 4e,f is the tidal EPF divergence (contours) that is computed
210 from the average daily cycle and includes all tidal components. The HIAMCM shows an
211 eastward tidal EPF divergence that maximizes around 0.003 hPa (85 km) and exceeds 15
212 $\text{ms}^{-1}\text{d}^{-1}$ in the southern summer mesosphere. Hence, the attenuation of the DE3 below

213 the summer mesopause gives rise to a significant contribution (10-20%) to the driving of
214 the equatorward residual circulation. The corresponding effect during July is very small.

4. Tidal components in the summer mesosphere from SABER

215 MLT temperatures are routinely measured by the Sounding the Atmosphere using
216 Broadband Emission Radiometry (SABER) instrument onboard the TIMED satellite
217 [Russell III *et al.*, 1999]. Standard tidal diagnostics of SABER have been detailed in
218 earlier papers [i.e., Forbes *et al.*, 2008] and requires combining 60 days of observations for
219 complete local solar time coverage. Furthermore, the spacecraft performs a yaw maneuver
220 approximately every 60 days (half of its precession period) to prevent SABER from point-
221 ing directly at the Sun. This changes the latitude coverage of the measurements from
222 55°S–85°N to 85°S–55°N, and vice versa. Yaws happened on 31 December 2016 and on 14
223 July 2006, and SABER was looking into the wrong hemisphere in January 2017 and late
224 July 2006. We therefore compare here observations for 21-30 December 2016 and 3-14 July
225 2006 to the model results for 1-20 January 2017 and 6-20 July 2006, respectively. To avoid
226 the 60-day averaging, we obtain a DE3 amplitude proxy as follows. For the 10-day periods
227 preceding the yaws, we fit zonal wave number 4 separately to observations made on the
228 ascending (asc) and descending (dsc) orbit nodes. A wave 4 observed in the satellite local
229 solar time frame of reference is, generally speaking, a superposition of a stationary wave 4
230 and various non-migrating tides (DW5, DE3, SW6, SE2, and the terdiurnal components
231 TW7 and TE1) [Oberheide *et al.*, 2011]. The local time difference between the asc and dsc
232 observations in the hemisphere of interest is about 14 hours. Differencing asc and dsc fits

233 thus amplifies the DE3 amplitudes (factor of 2) while minimizing semidiurnal, terdiurnal,
234 and stationary wave signals.

235 Figure 5 shows the results for December 2016 and July 2006. The patterns are struc-
236 turally similar to the HIAMCM and MERRA-2 results (Figs. 2e,f and 3e,f). This includes
237 a stronger DE3 in the low-latitude MLT during July. In particular, SABER shows a
238 pronounced middle to high-latitude DE3 maximum in the summer mesosphere during
239 December that tilts towards higher altitudes with increasing latitude. A maximum DE3
240 amplitude of ~ 3 K is found at ~ 75 km and 50° S. The middle to high-latitude DE3 in
241 July from SABER is less pronounced, which is also consistent with the model result. Note
242 that the high-latitude SABER maximum for December does not extend much above 80
243 km. Whether this difference with respect to the HIAMCM is due to some interference in
244 the asc-dsc differences or other effects cannot be resolved with the data at hand.

5. Conclusions

245 We have documented the presence of an unexpected DE3 tide in the southern summer
246 mesosphere at middle to high latitudes. We first found this DE3 in a simulation of Jan-
247 uary 2017 using a GW-resolving GCM (HIAMCM). We showed that the model result is
248 consistent with MERRA-2 reanalysis and a new tidal analysis of SABER temperature
249 data. Moreover, the large diurnal tidal amplitude from the DE3 is quantitatively con-
250 sistent with previous lidar measurements at Antarctica [Lübken *et al.*, 2011]. From the
251 zonal-mean analysis we concluded that the attenuation of the DE3 below the summer
252 mesopause gives a significant eastward EPF divergence that contributes about 10-20% to
253 the driving of the equatorward residual circulation. We also analyzed a period during

254 July 2006 and found that the diurnal tide in the northern summer mesosphere is mainly
255 a combination of eastward non-migrating tides (DE1, DE2, and DE3). The overall diur-
256 nal tide is weaker than in the southern summer mesosphere, which is in agreement with
257 ground-based measurements by *Gerding et al.* [2013].

258 A strong DE3 in the southern summer mesosphere is usually not found in linear tidal
259 models. We also inspected data from a GCM with parameterized GWs [*Becker, 2017*] and
260 found no indication of DE components in the mesosphere at middle to high latitudes (not
261 shown in this paper). Note that all these models exclude important aspects of GW-tidal
262 interactions [e.g. *Senf and Achatz, 2011*]. This suggests that the unexpected DE3 in the
263 southern summer mesosphere is simulated only in models with resolved GWs.

264 We analyzed only the northern winter 2016-2017 to document the DE3 in the south-
265 ern summer mesosphere. Analyses of other periods are necessary to determine whether
266 our results apply more generally. Also, a detailed investigation of the GW-tidal interac-
267 tions and other possible mechanisms that may explain the DE3 in the southern summer
268 mesosphere and hemispheric differences of tidal components is demanded by our findings.
269 These efforts are, however, beyond the scope of this paper and will be subject to future
270 studies.

271 **Acknowledgments.** EB acknowledges support by NASA grants 80NSSC19K0834
272 and 80NSSC21M0180. JO was supported by NASA grants 80NSSC19K0258 and
273 80NSSC20K1353. We thank the Leibniz Institute of Atmospheric Physics at the Uni-
274 versity of Rostock and the NASA High-End Computing Capability (HECC) for providing

275 HPC facilities used in this study. The comments by two anonymous reviewers are greatly
276 acknowledged.

277 **Open Research.** Model documentations can be found in in *Becker and Vadas* [2020],
278 *Becker et al.* [2022a], and *Becker et al.* [2022b]. Model data shown in this study can be
279 downloaded from NWRA’s website under

280 <https://www.cora.nwra.com/~erich.becker/Becker-Oberheide-GRL-2023-files>.

281 The MERRA-2 reanalysis data are publicly available at

282 <https://goldsmr5.gesdisc.eosdis.nasa.gov/data/MERRA2/M2I3NVASM.5.12.4/2017>.

283 SABER v2.07 data can be downloaded from

284 <https://saber.gats-inc.com/data.php>.

References

285 Achatz, U., N. Grieger, and H. Schmidt (2008), Mechanisms controlling the diurnal solar
286 tide: Analysis using a gcm and a linear model, *J. Geophys. Res.*, *113*(A08303), doi:
287 10.1029/2007JA012967.

288 Akmaev, R. A. (2001), Simulation of large-scale dynamics in the mesosphere and lower
289 thermosphere with the doppler-spread parameterization of gravity waves 2. eddy mixing
290 and the diurnal tide, *J. Geophys. Res.*, *106*(D1), 1205–1213, doi:10.1029/2000JD900520.

291 Becker, E. (2017), Mean-flow effects of thermal tides in the mesosphere and lower ther-
292 mosphere, *J. Atmos. Sci.*, *74*, 2043–2063, doi:10.1175/JAS-D-16-0194.1.

293 Becker, E., and S. L. Vadas (2018), Secondary gravity waves in the winter mesosphere:
294 Results from a high-resolution global circulation model, *J. Geophys. Res. Atmos.*, *123*,

- 295 doi:10.1002/2017JD027460.
- 296 Becker, E., and S. L. Vadas (2020), Explicit global simulation of gravity waves in the
297 thermosphere, *J. Geophys. Res. Space Phys.*, doi:10.1029/2020JA028034.
- 298 Becker, E., S. L. Vadas, K. Bossert, V. L. Harvey, C. Zülicke, and L. Hoffmann (2022a),
299 A high-resolution whole-atmosphere model with resolved gravity waves and specified
300 large-scale dynamics in the troposphere and stratosphere, *J. Geophys. Res. Atmos.*,
301 *127*(e2021JD035018), doi:10.1029/2021JD035018.
- 302 Becker, E., L. Goncharenko, V. L. Harvey, and S. L. Vadas (2022b), Multi-step vertical
303 coupling during the january 2017 sudden stratospheric warming, *J. Geophys. Res. Space*
304 *Phys.*, *127*(e2022JA030866), doi:10.1029/2022JA030866.
- 305 Bosilovich, M. G., S. Akella, L. Coy, R. Cullather, C. Draper, R. Gelaro, R. Kovach,
306 Q. Liu, A. Molod, P. Norris, K. Wargan, W. Chao, R. Reichle, L. Takacs, Y. Vikhliayev,
307 S. Bloom, A. Collopy, S. Firth, G. Labow, G. Partyka, S. Pawson, O. Reale, S. D. Schu-
308 bert, and M. Suarez (2015), MERRA-2: Initial evaluation of the climate, *NASA Tech.*
309 *Rep. Series on Global Modeling and Data Assimilation*, *NASA/TM-2015-104606/Vol.*
310 *43*, Goddard Space Flight Center.
- 311 Chapman, S., and R. S. Lindzen (1969), Atmospheric tides, *Space Sci. Rev.*, *10*(1), 3–188,
312 doi:10.1007/BF00171584.
- 313 Fomichev, V. I., W. E. Ward, S. R. Beagley, C. McLandress, J. C. McConnell, N. A.
314 McFarlane, and T. G. Shepherd (2002), Extended Canadian Middle Atmosphere Model:
315 Zonal-mean climatology and physical parameterizations, *J. Geophys. Res.*, *107*(D10),
316 ACL9–1–ACL9–14, doi:10.1029/2001JD000479.

- 317 Forbes, J. M. (1984), Middle atmosphere tides, *J. Atmos. Terr. Phys.*, *46*(11), 1049–1067,
318 doi:10.1016/0021-9169(84)90008-4.
- 319 Forbes, J. M., X. Zhang, S. Palo, J. Russell, C. J. Mertens, , and M. Mlynczak (2008),
320 Tidal variability in the ionospheric dynamo region, *J. Geophys. Res. Space Phys.*,
321 *113*(A02310), doi:10.1029/2007JA012737.
- 322 Gelaro, R., W. McCarty, M. J. Suárez, R. Todling, A. Molod, L. Takacs, C. A. Ran-
323 dles, A. Darmenov, M. G. Bosilovich, R. Reichle, K. Wargan, L. Coy, R. Cullather,
324 C. Draper, S. Akella, V. Buchard, A. Conaty, A. M. Da Silva, W. Gu, G.-K. Kim,
325 R. Koster, R. Lucchesi, D. Merkova, J. E. Nielsen, G. Partyka, S. Pawson, W. Putman,
326 M. Rienecker, S. D. Schubert, M. Sienkiewicz, and B. Zhao (2017), The Modern-Era
327 Retrospective analysis for Research and Applications, Version 2 (MERRA-2), *J. Clim.*,
328 *30*, 5419–5453, doi:10.1175/JCLI-D-16-0758.1.
- 329 Gerding, M., M. Kopp, P. Hoffmann, J. Höffner, and F.-J. Lübken (2013), Diur-
330 nal variations of midlatitude NLC parameters observed by daylight-capable lidar
331 and their relation to ambient parameters, *Geophys. Res. Lett.*, *40*, 6390–6394, doi:
332 10.1002/2013GL057955.
- 333 Hagan, M. E., and J. M. Forbes (2002), Migrating and nonmigrating diurnal tides in the
334 middle and upper atmosphere excited by tropospheric latent heat release, *J. Geophys.*
335 *Res.*, *107*(D24), doi:10.1029/2001JD001236.
- 336 Harvey, V. L., N. Pedatella, E. Becker, and C. E. Randall (2022), Evaluation of po-
337 lar winter mesopause wind in WACCMX+DART, *J. Geophys. Res. Atmos.*, doi:
338 10.1029/2022JD037063.

- 339 Karlsson, B., and E. Becker (2016), How does interhemispheric coupling contribute to
340 cool down the summer polar mesosphere?, *J. Clim.*, *29*, 8807–8821, doi:10.1175/JCLI-
341 D-16-0231.1.
- 342 Kishore Kumar, G., W. Singer, J. Oberheide, N. Grieger, P. P. Batista, D. M.
343 Riggin, H. Schmidt, and B. R. Clemesha (2014), Diurnal tides at low latitudes:
344 Radar, satellite, and model results, *J. Atmos. Sol. Terr. Phys.*, *118*, 96–105, doi:
345 10.1016/j.jastp.2013.07.005.
- 346 Körnich, K., and E. Becker (2010), A simple model for the interhemispheric cou-
347 pling of the middle atmosphere circulation, *Adv. Space Res.*, *45*, 661–668, doi:
348 10.1016/j.asr.2009.11.001.
- 349 Kumari, K., J. Oberheide, and X. Lu (2020), The tidal response in the meosphere/lower
350 thermosphere to the Madden-Julian oscillation observed by SABER, *Geophys. Res.*
351 *Lett.*, *47*(e2020GL089172), doi:10.1029/2020GL089172.
- 352 Lieberman, R. S., J. Oberheide, M. E. Hagan, E. E. Remsberg, and L. L. Gordley (2014),
353 Variability of diurnal tides and planetary waves during november 1978–may 1979, *J. .*
354 *Atmos. Sol. Terr. Phys.*, *66*, 571–528, doi:10.1016/j.jastp.2004.01.006.
- 355 Liu, H.-L., B. T. Foster, M. E. Hagan, J. M. McInerney, A. Maute, L. Qian, A. D.
356 Richmond, R. G. Roble, S. C. Solomon, R. R. Garcia, D. Kinnison, D. R. Marsh,
357 A. K. Smith, J. Richter, F. Sassi, and J. Oberheide (2010), Thermosphere extension
358 of the whole atmosphere community climate model, *J. Geophys. Res.*, *115*(A12301),
359 doi:10.1029/2010JA015586.

- 360 Lübken, F.-J., J. Höffner, T. P. Viehl, B. Kaifler, and R. J. Morris (2011), First measure-
361 ments of thermal tides in the summer mesopause region at Antarctic latitudes, *Geophys.*
362 *Res. Lett.*, *38*, doi:10.1029/2011GL050045.
- 363 Oberheide, J., and J. M. Forbes (2008), Tidal propagation of deep tropical cloud signatures
364 into the thermosphere from TIMED observations, *Geophys. Res. Lett.*, *35*(L04816114),
365 doi:10.1029/2007GL032397.
- 366 Oberheide, J., J. M. Forbes, K. Häusler, Q. Wu, and S. L. Bruinsma (2009), Tropospheric
367 tides from 80 to 400 km: Propagation, interannual variability, and solar cycle effects,
368 *J. Geophys. Res.*, *114*(D00I05), doi:10.1029/2009JD012388.
- 369 Oberheide, J., J. M. Forbes, X. Zhang, and S. L. Bruinsma (2011), Wave-driven variability
370 in the ionosphere-thermosphere-mesosphere system from timed observations: What con-
371 tributes to the “wave 4”?, *J. Geophys. Res.*, *116*(A01306), doi:10.1029/2010JA015911.
- 372 Pedatella, N. M., J. Oberheide, E. K. S. a nd H.-L. Liu, J. L. Anderson, and K. Reader
373 (2016), Short-term nonmigrating tide variability in the mesosphere, thermosphere, and
374 ionosphere, *J. Geophys. Res. Space Phys.*, *121*, 3621–3633, doi:10.1002/2016JA022528.
- 375 Russell III, J. M., M. G. Mlynczak, L. L. Gordley, J. J. Tansock Jr., and R. W. Esplin
376 (1999), Overview of the saber experiment and preliminary calibration results, *Proc.*
377 *spie 3756*, Optical Spectroscopic Techniques and Instrumentation for Atmospheric and
378 Space Research III (20 October 1999), doi:10.1117/12.366382.
- 379 Sakazaki, T., M. Fujiwara, X. Zhang, M. E. Hagan, and J. M. Forbes (2012), Diurnal
380 tides from the troposphere to the lower mesosphere as deduced from TIMED/SABER
381 satellite data and six global reanalysis data sets, *J. Geophys. Res.*, *117*(D13108), doi:

382 10.1029/2011JD017117.

383 Senf, F., and U. Achatz (2011), On the impact of middle-atmosphere thermal tides
384 on the propagation and dissipation of gravity waves, *J. Geophys. Res.*, *116*, doi:
385 10.1029/2011JD015794.

386 Smith, A. K. (2012), Global dynamics of the MLT, *Surv. Geophys.*, *33*, 1177–1230, doi:
387 10.1007/s10712-012-9196-9.

388 Smith, A. K., R. R. Garcia, D. R. Marsh, and J. H. Richter (2011), WACCM simulations of
389 the mean circulation and trace species transport in the winter mesosphere, *J. Geophys.*
390 *Res.*, *116*(D20115), doi:10.1029/2011JD016083.

391 Vadas, S. L., and E. Becker (2019), Numerical modeling of the generation of tertiary
392 gravity waves in the mesosphere and thermosphere during strong mountain wave events
393 over the Southern Andes, *J. Geophys. Res. Space Phys.*, doi:10.1029/2019JA026694.

394 Vitharana, A., X. Zhu, J. Du, J. Oberheide, and W. Ward (2019), Statistical modeling
395 of tidal weather in the mesosphere and lower thermosphere, *J. Geophys. Res. Atmos.*,
396 *124*, 9011–9027, doi:10.1029/2019JD030573.

397 Ward, W. E., J. Oberheide, L. P. Goncharenko, T. Nakamura, P. Hoffmann, W. Singer,
398 L. C. Chang, J. Du, D.-Y. Wang, P. Batista, B. Clemesha, A. H. Manson, C. Meek, D. M.
399 Riggan, C.-Y. She, T. Tsuda, and T. Yuan (2010), On the consistency of model, ground-
400 based and satellite observations of tidal signatures: Initial results from the cawses tidal
401 campaigns, *J. Geophys. Res.*, *115*(D07107), doi:10.1029/2009JD012593.

402 Zhang, X., J. M. Forbes, and M. E. Hagan (2010a), Longitudinal variation of tides in
403 the mlt region: 1. Tides driven by troposphere net radiative heating, *J. Geophys. Res.*,

404 115(A06316), doi:10.1029/2009JA014897.

405 Zhang, X., J. M. Forbes, and M. E. Hagan (2010b), Longitudinal variation of tides in the

406 mlt region: 2. Relative effects of solar radiative and latent heating, *J. Geophys. Res.*,

407 115(A06317), doi:10.1029/2009JA014898.

408 Zhang, X., J. M. Forbes, and M. E. Hagan (2012), Seasonal-latitudinal variation of the

409 eastward-propagating diurnal tide with zonal wavenumber 3 in the MLT: Influences of

410 heating and background wind distribution, *J. Atmos. Sol.-Terr. Phys.*, 78-79, 37-43,

411 doi:10.1016/j.jastp.2011.03.005.

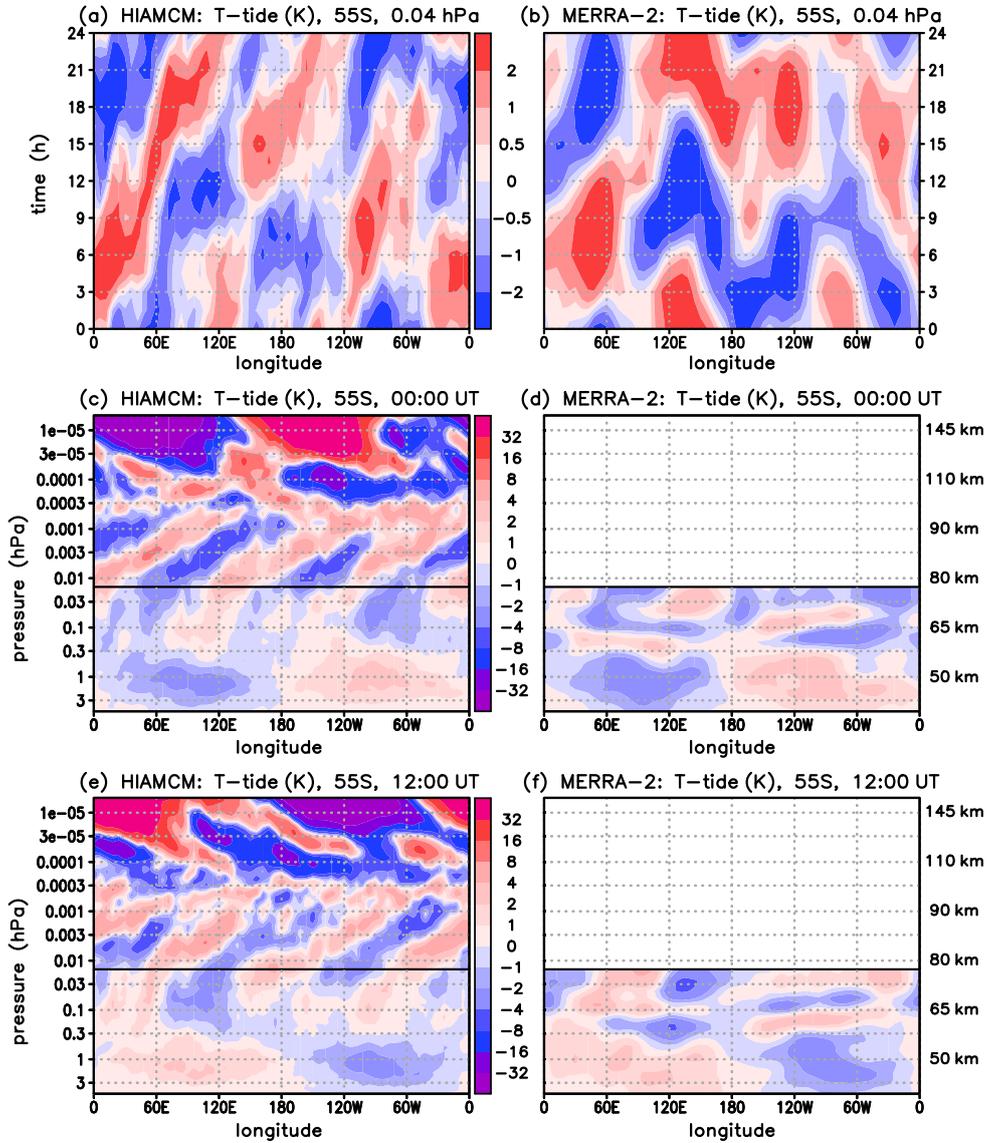


Figure 1. Temperature tide (K) at 55°S for 1-20 January 2017 from the HIAMCM (left column) and from MERRA-2 (right column). First row: Longitude-time plots at 0.04 hPa. Second (third) row: Longitude-height plots at 0 UT (12 UT). Approximate heights are given on the right-hand sides of panels d and f.

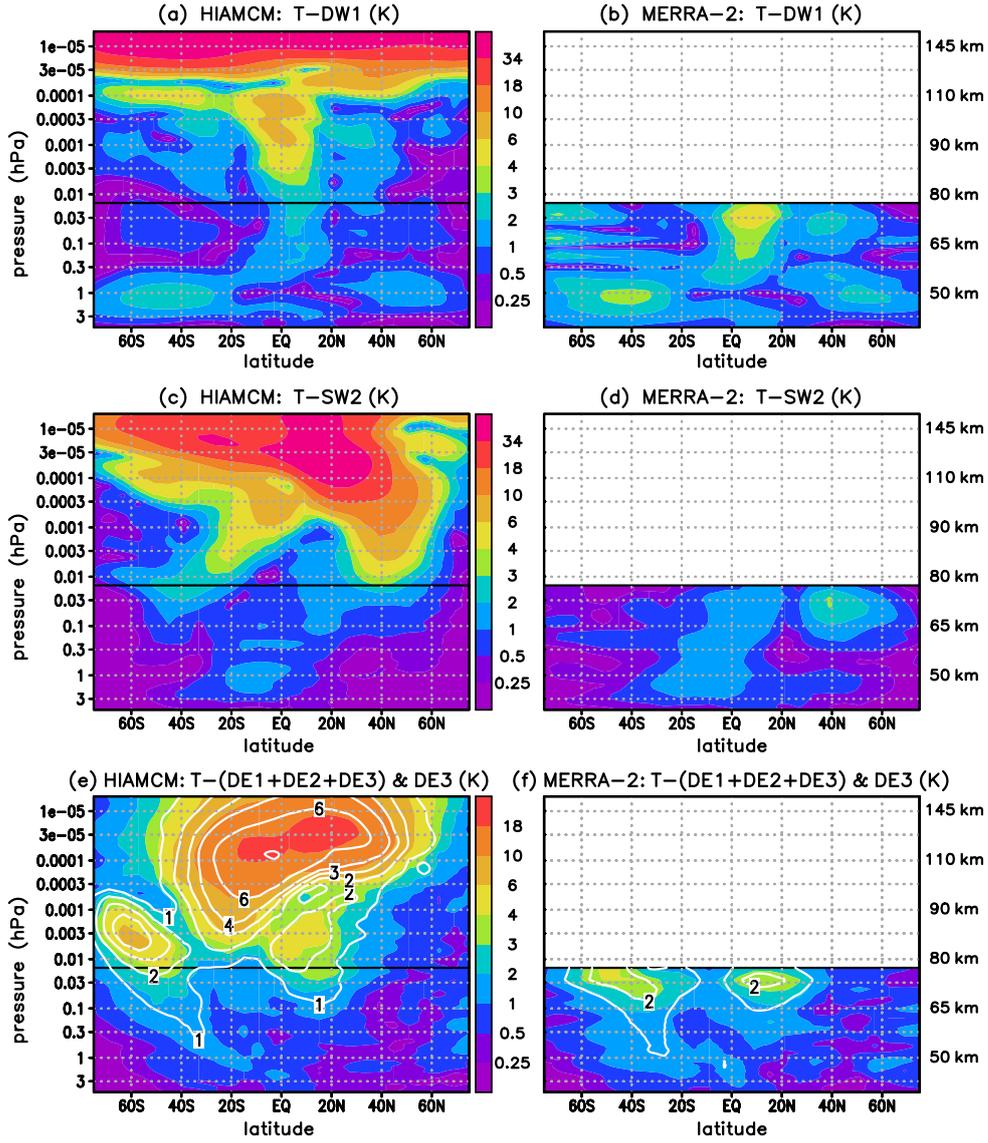


Figure 2. Tidal temperature amplitudes (K) for the 1-20 January 2017 period from the HIAMCM (left column) and from MERRA-2 (right column). First row: DW1. Second row: SW2. Third row: Sum of eastward propagating diurnal tides with zonal wavenumbers $s = 1$ to 3 (DE1+DE2+DE3, colours) and amplitude of the DE3 (white contours for 1, 2, 3, 4, 6, 8 K). Approximate heights are given on the right-hand sides of panels b, d, and f.

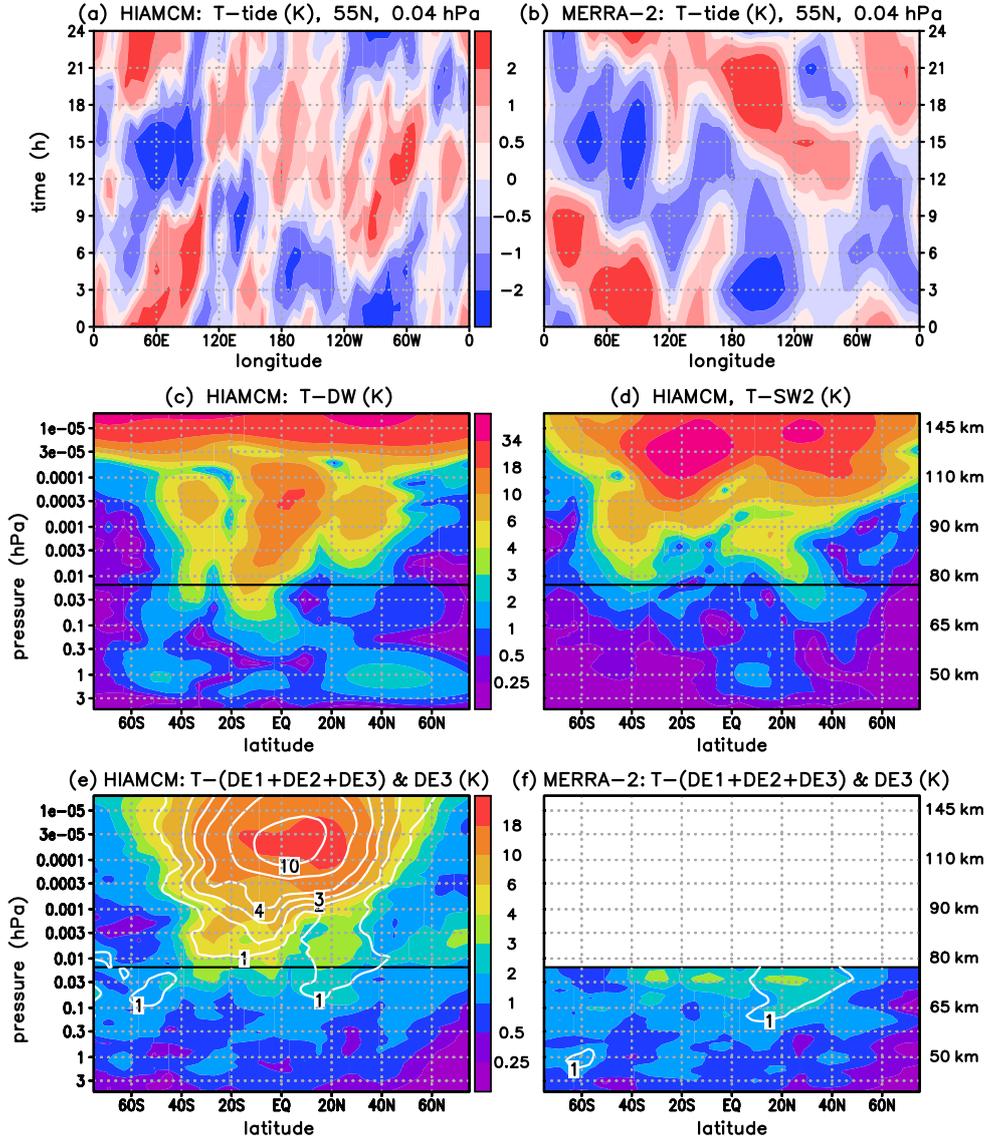


Figure 3. (a),(b) Same as Fig. 1a,b, but for 6-20 July 2006 and at 55°N. (c),(d) Same Fig. 2a,c, but for 6-20 July 2006. (e),(f) Same Fig. 2e,f, but for 6-20 July 2006.

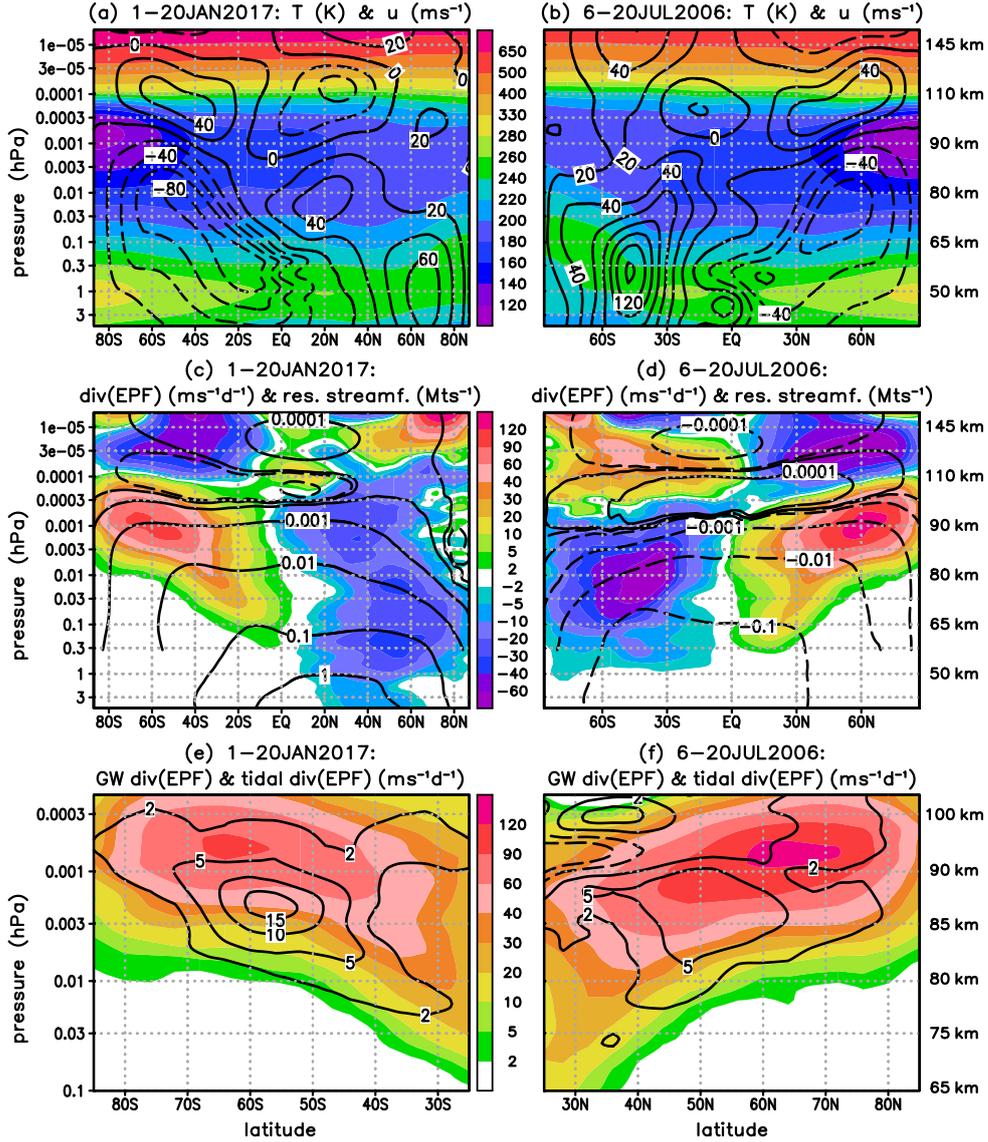


Figure 4. Zonal-mean circulation and wave driving from the HIAMCM for 1-20 January 2017 (left) and 6-20 July 2006 (right). First row: Temperature (colours) and zonal wind (contour interval 20 ms^{-1}). Second row: EPF divergence (colors, unit $\text{ms}^{-1}\text{d}^{-1}$) and residual mass streamfunction (contours for $\pm 10^{-4}$, $\pm 10^{-3}$, $+0.01$, $+0.1$, $+1 \text{ Mt s}^{-1}$ in (c) and for $\pm 10^{-4}$, $\pm 10^{-3}$, -0.01 , -0.1 Mt s^{-1} in (d), $1 \text{ Mt} = 10^9 \text{ kg}$). Third row: EPF divergence due to GWs (colors, unit $\text{ms}^{-1}\text{d}^{-1}$) and tides (contours for ± 2 , ± 5 , 10 , $15 \text{ ms}^{-1}\text{d}^{-1}$) in the summer mesosphere.

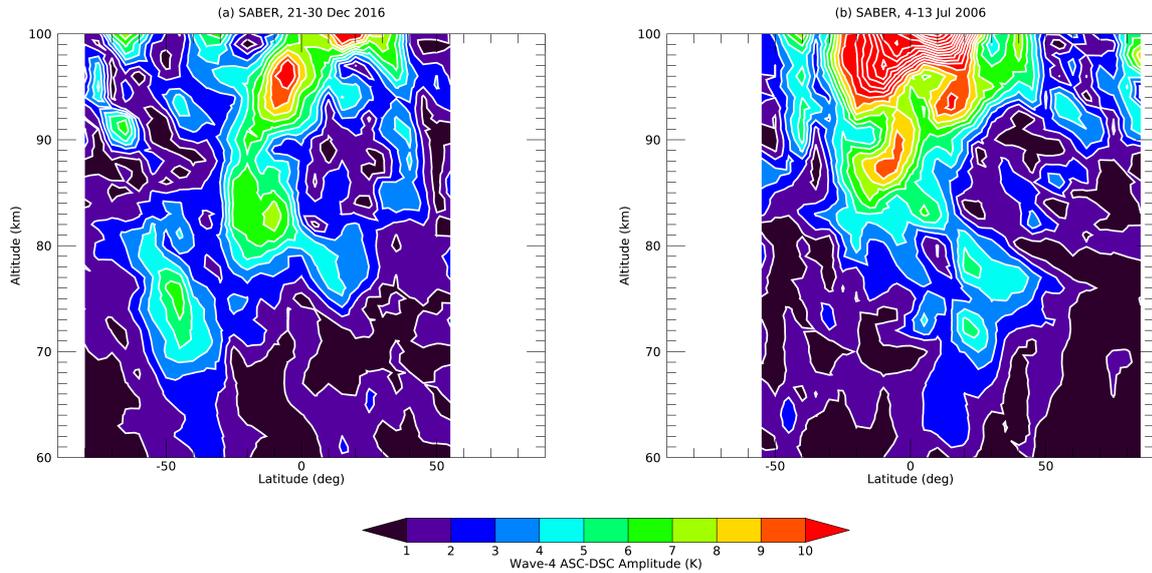


Figure 5. DE3 temperature amplitude estimates from SABER for (a) 21-30 December 2016 and (b) 3-13 July 2006: SABER wave-4 asc-dsc difference amplitudes. The amplitudes need to be divided by a factor of 2 for comparison with Figs. 2e,f and 3e,f.